On the muon scale of air showers and its application to the AGASA data

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Muon scale and its application to AGASA data

Motivation and objective

Evidence of a muon deficit in air shower simulations [1, 2, 3, 4].



Objectives

 Define two estimators of z, study their properties & syst. Compute their values from AGASA data.

$z_{\ln\langle \cdot angle}$ and $z_{\langle \ln \cdot angle}$

We introduce the average over an energy bin in two ways:

Experiments with different resolution: same $\ln \langle N_{\mu,\text{data}}^{\text{det}} \rangle \rightarrow \text{different } \langle \ln N_{\mu,\text{data}}^{\text{det}} \rangle$

 $\rightarrow z_{\ln(\cdot)}$ is better for comparing different experiments

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4 3 5 4 5

Sources of systematics	Syst. in $z_{\ln\langle \cdot \rangle}$	Syst. in $z_{\langle \ln \cdot \rangle}$
Composition bias in data $\langle N_{\mu,\rm data}^{\rm det}\rangle$	\checkmark	\checkmark
Mismodeled detector effects in the mean $\langle N_{\mu,\{\rm p,Fe\}}^{\rm det}\rangle$	\checkmark	\checkmark
Mismodeled detector resolution		\checkmark

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Mismodeled detector resolution		\checkmark

Syst. in $z_{\langle \ln \cdot \rangle}$ from mismodeled detector resolution

If we assume

- No mismodeled detector effects in the mean $\langle N_{\mu,\{p,Fe\}}^{det} \rangle$ (already in syst.).
- $\langle \ln N_{\mu} \rangle \approx \ln \langle N_{\mu} \rangle \frac{1}{2} \left(\text{RSD}_{\text{tot}}[N_{\mu}] \right)^2$ (Eq. (3)).
- Heitler-Matthews model (see Eq. (4))

Then...

Syst. in $z_{\langle \ln \cdot \rangle}$ from mismodeled detector resolution



Fig. 2: Syst. err. in $z_{(\ln \cdot)}$ as a function of the true det. res. (x-axis) and of the difference between the mismodeled and true det. res. (y-axis).

Syst. in $z_{\langle \ln \cdot \rangle}$ from mismodeled detector resolution



Fig. 3: E.g.: The total det. resolution of AGASA is ~ 50 % in the best case. Syst. err of ± 0.07 in $z_{\langle \ln \cdot \rangle}$ is attained when the resolution is mismodeled in ± 6 %.

Muon deficit scale

To compute a muon deficit scale, we need a reference. Heitler-Matthews model [6]

$$\ln N_{\mu} = (1 - \beta) \ln A + \beta \ln(E/\xi_c),$$
power-law index $\beta \approx 0.9$ critical energy constant

In the Heitler-Matthews model:
$$z_{\langle \ln \cdot \rangle \max}^{\text{HM}} = \frac{\langle \ln A \rangle}{\ln 56}$$
.

using mass composition model

We therefore compute:

$$\Delta z_{\ln\langle\cdot\rangle} = z_{\ln\langle\cdot\rangle} - z_{\langle\ln\cdot\rangle\,\max}^{\rm HM}$$

$$\Delta z_{\langle\ln\cdot\rangle} = z_{\langle\ln\cdot\rangle} - z_{\langle\ln\cdot\rangle\,\max}^{\rm HM}$$
(5)
(6)

primary energy

(4)

Sources of systematics	Syst. in $\Delta z_{\ln\langle \cdot \rangle}$	Syst. in $\Delta z_{\langle \ln \cdot \rangle}$
Propagated from $z_{\ln\langle\cdot\rangle}$ or $z_{\langle\ln\cdot\rangle}$	\checkmark	\checkmark
Syst. in the composition model in $z^{\rm HM}_{\langle \ln \cdot \rangle \rm mass}$	\checkmark	\checkmark
Deviations from the H.M. model in $z^{\rm HM}_{\langle \ln \cdot \rangle \rm mass}$	\checkmark	\checkmark
Bias from shower-to-shower fluctuations	\checkmark	

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Syst. in the composition model in $z^{\rm HM}_{\langle \ln \cdot \rangle \rm mass}$	\checkmark	\checkmark
Deviations from the H.M. model in $z^{\rm HM}_{\langle \ln \cdot \rangle \rm mass}$	\checkmark	\checkmark
Bias from shower-to-shower fluctuations	\checkmark	

Bias in $\Delta z_{\ln(\cdot)}$ from sh-sh fluctuations in $z_{\langle \ln \cdot \rangle \text{ mass}}^{\text{HM}}$

- If $\Delta z_{\ln\langle\cdot\rangle}$ were computed taking $z_{\ln\langle\cdot\rangle}$ mass as reference, there would be no bias.
- However, we take as a reference $z_{(\ln \cdot) \text{ mass}}^{\text{HM}}$, the predicted value of $z_{(\ln \cdot) \text{ mass}}$.

∠det. sim. using mass composition model

$$z_{\langle \ln \cdot \rangle \operatorname{mass}} = \frac{\langle \ln N_{\mu, \operatorname{mass}}^{\operatorname{det}} \rangle - \langle \ln N_{\mu, p}^{\operatorname{det}} \rangle}{\langle \ln N_{\mu, Fe}^{\operatorname{det}} \rangle - \langle \ln N_{\mu, p}^{\operatorname{det}} \rangle}$$

$$\approx \frac{\left[\ln \langle N_{\mu, \operatorname{mass}}^{\operatorname{det}} \rangle - \ln \langle N_{\mu, p}^{\operatorname{det}} \rangle - \frac{1}{2} \left[(\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, \operatorname{mass}}])^2 - (\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, p}])^2 \right]}{\ln \langle N_{\mu, Fe}^{\operatorname{det}} \rangle - \ln \langle N_{\mu, p}^{\operatorname{det}} \rangle} - \frac{1}{2} \left[(\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, \operatorname{Fe}}])^2 - (\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, p}])^2 \right]}{\langle n_{\lambda, Fe} \rangle - \ln \langle N_{\mu, p}^{\operatorname{det}} \rangle} - \frac{1}{2} \left[(\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, Fe}])^2 - (\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, p}])^2 \right]}{\langle n_{\lambda, Fe} \rangle - \ln \langle N_{\mu, p}^{\operatorname{det}} \rangle} - \frac{1}{2} \left[(\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, Fe}])^2 - (\operatorname{RSD}_{\operatorname{sh-sh}}[N_{\mu, p}])^2 \right]} \right]}{\langle n_{\lambda, Fe} \rangle - \ln \langle N_{\mu, p}^{\operatorname{det}} \rangle}$$

$$(7)$$

 $\rightarrow z_{\langle \ln \cdot \rangle \text{ mass}}$ depends on sh-sh fluctuations, while $z_{\ln \langle \cdot \rangle \text{ mass}}$ does not.

Worst case scenario: ~ 50 % p + 50 % Fe $\rightarrow |z_{\ln\langle\cdot\rangle mass} - z_{\langle\ln\cdot\rangle mass}| \leq 0.07$ [5]. Furthermore, this bias in $\Delta z_{\ln\langle\cdot\rangle}$ can be corrected.

Overall systematics comparison



Fig. 4: In typical scenarios, it is better to use $\Delta z_{\ln(\cdot)}$.

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 10^{1}

 $\rho_{\mu}\,[\mathrm{m}^{-2}]$

 10^{-1}

 10^{-2} ė 00 0

Akeno Giant Air Shower Array



Fig. 5: Image extracted from Ref. [7].



19.4 19.6 19.8 20.0 20.2

No detected muons

Data

8

o AGASA

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100 0



Sim. library described in Ref. [1]. Det. effects are accounted for analytically.

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Results



Fig. 7:
$$z_{\ln\langle\cdot\rangle} = \frac{\ln\langle\rho_{\mu,data}^{det}\rangle - \ln\langle\rho_{\mu,p}\rangle}{\ln\langle\rho_{\mu,Fe}\rangle - \ln\langle\rho_{\mu,p}\rangle}$$
 vs $z_{\langle\ln\cdot\rangle} \approx \frac{\langle\ln\rho_{\mu,data}^{det}\rangle - \ln\langle\rho_{\mu,p}\rangle + \frac{1}{2} [(\text{RSD}_{\text{sh-sh}}[\rho_{\mu,p}])^2 + (\text{RSD}_{\text{not sh-sh}}[\rho_{\mu}])^2]}{\ln\langle\rho_{\mu,Fe}\rangle - \ln\langle\rho_{\mu,p}\rangle + \frac{1}{2} [(\text{RSD}_{\text{sh-sh}}[\rho_{\mu,p}])^2 - (\text{RSD}_{\text{sh-sh}}[\rho_{\mu,Fe}])^2]}{(n \langle\rho_{\mu,Fe}\rangle - \ln\langle\rho_{\mu,p}\rangle + \frac{1}{2} [(\text{RSD}_{\text{sh-sh}}[\rho_{\mu,p}])^2 - (\text{RSD}_{\text{sh-sh}}[\rho_{\mu,Fe}])^2]}$

Results

Fig. 8: Image adaptd from Ref. [4].



There is an agreement with Pierre Auger and Yakutsk array data. AGASA data support a muon deficit in air shower simulations.

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Muon scale and its application to AGASA data

Summary

Muon scale: $z_{\ln\langle\cdot\rangle}$ vs $z_{\langle\ln\cdot\rangle}$

- $z_{(\ln \cdot)}$ suffers from systematics if the detector resolution is mismodeled.
- We provided a way to estimate this systematic for any experiment.

 $z_{\ln\langle\cdot\rangle}$ is always better for comparing results from different experiments.

Muon deficit scale: $\Delta z_{\ln\langle \cdot \rangle}$ vs $\Delta z_{\langle \ln \cdot \rangle}$

- $-\Delta z_{\ln\langle\cdot\rangle}$ has a bias from shower-to-shower fluctuations in $z_{\ln\langle\cdot\rangle}^{\text{HM}} = \langle \ln A \rangle / \ln 56$.
- The bias in $\Delta z_{\ln\langle\cdot\rangle}$ is typically smaller than the syst. in $\Delta z_{\langle\ln\cdot\rangle}$ + correctible. $\Delta z_{\ln\langle\cdot\rangle}$ is generally better for estimating the deficit.

\boldsymbol{z} from AGASA data

They are in agreement with Pierre Auger and Yakutsk Array data.
 AGASA data constitute further evidence of a muon deficit.

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