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A two-zone emission model for Blazars and the role of Accretion Disk MHD winds

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Ευρωπαϊκή Ένωση
European Social Fund

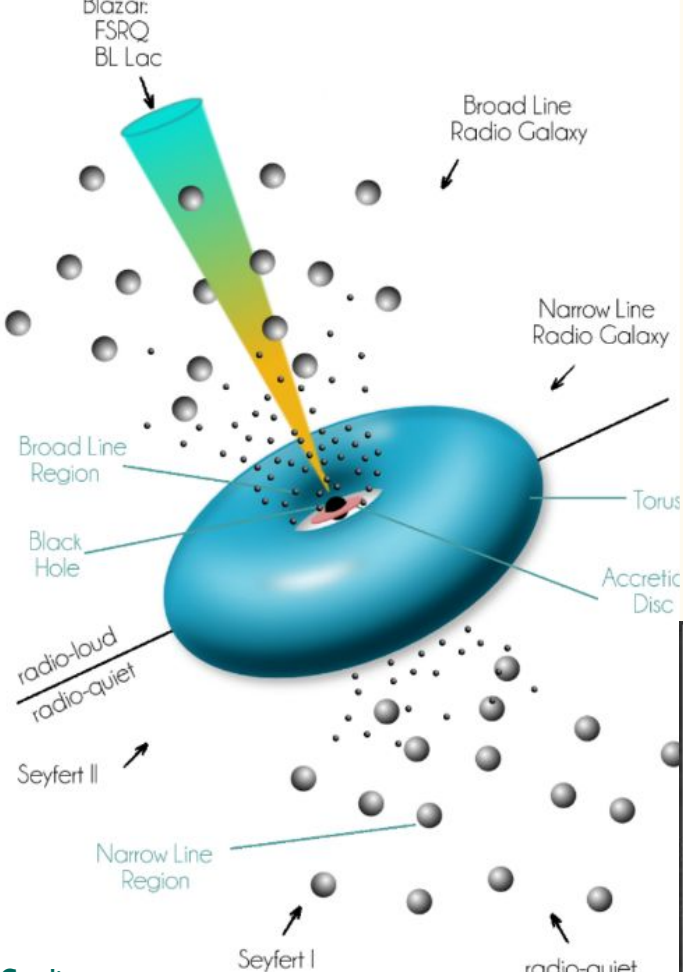
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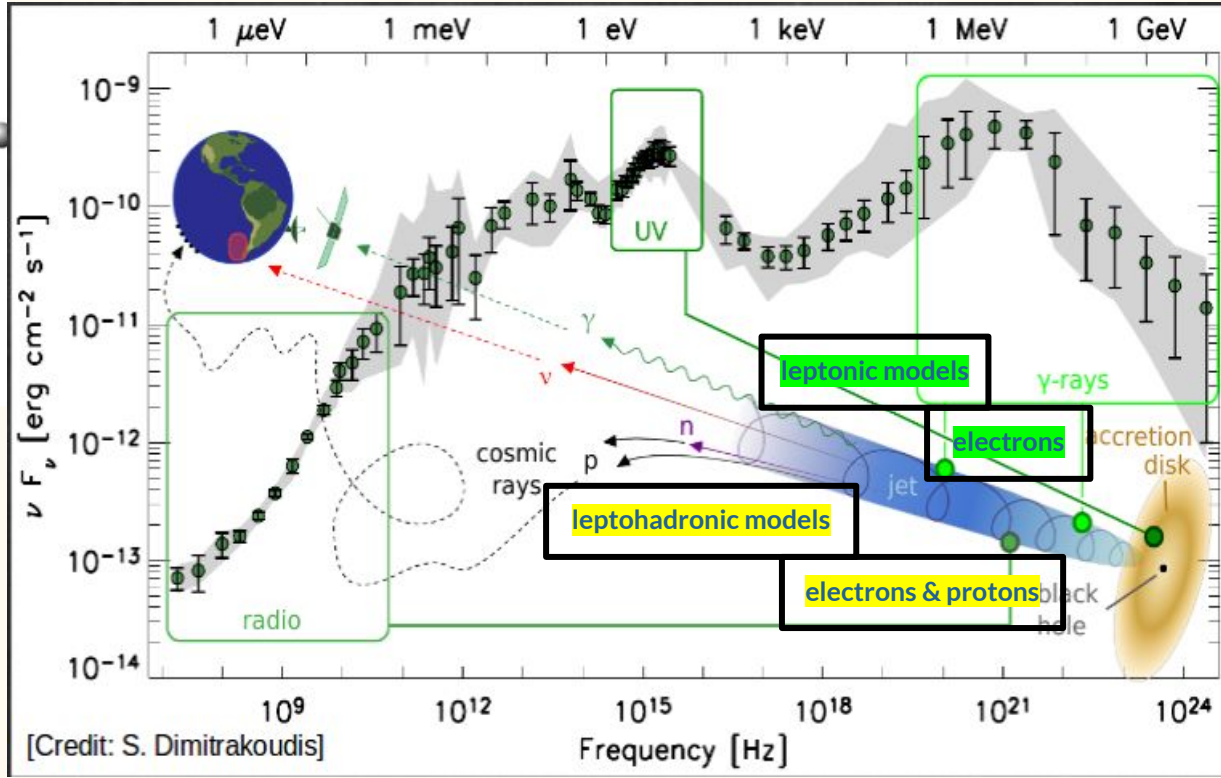


ανάπτυξη - εργασία - αλληλεγγύη

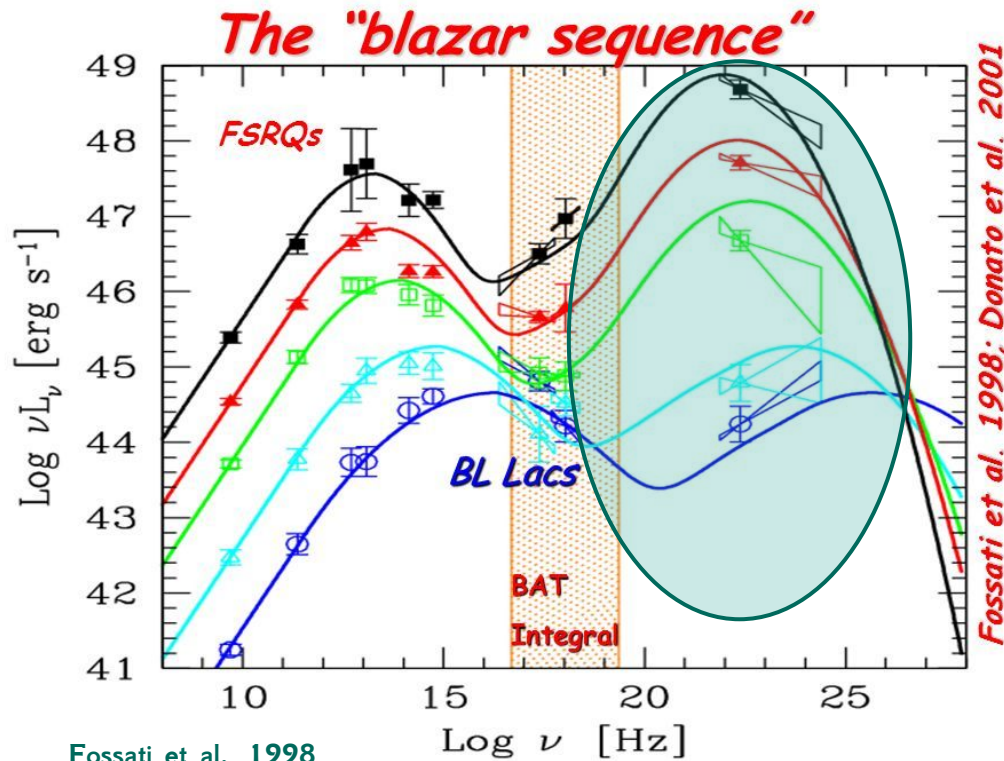




Credit:
<http://www.sternwarte.uni-erlangen.de/krauss>



Target Photon Fields



- Accretion Disk Photons (Dermer et al., 1992, Dermer and Schlickeiser, 1993 ++)
- Broad Line Region (Sikora et al., 1994, Blandford and Levinson, 1995, Ghisellini and Madau, 1996, Dermer et al., 1997, Finke, 2013 ++)
- Photons from torus (Blazejowski et al., 2000)
- Synchrotron emission from other regions of the jet (Georganopoulos and Kazanas, 2003, Ghisellini and Tavecchio, 2008)
- Photons which are scattered on Accretion Disk Wind particles (Boula et al., 2019)
- Synchrotron Photons (Marscher and Gear, 1985, Maraschi et al., 1992, Bloom and Marscher, 1996 ++)

Theoretical Emission Model

Basic parameters of a Leptonic Model

- Magnetic field strength
- Electrons luminosity
- Electrons distribution
- Energy density of the external photon field
- Bulk Lorentz factor $\Gamma = (1 - \beta^2)^{-1/2}$
- Doppler factor $\delta = [\Gamma(1 - \beta \cos \theta)]^{-1}$

**Related to the
mass accretion rate**

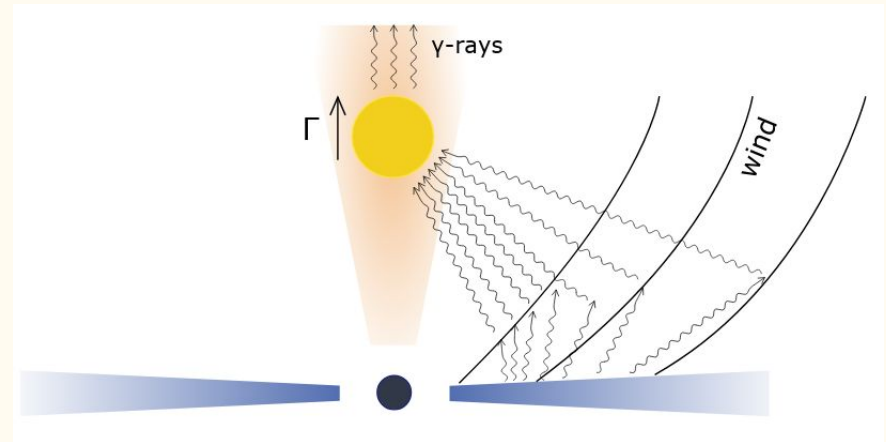


Image credit: S. Dimitrakoudis

External Photon Field

$$n(r, \theta) = n_0 (r_s / r)^p e^{5(\theta - \pi/2)} \quad n_0 = \frac{\eta_w \dot{m}}{2\sigma_T r_s}$$

$$n(r, \theta) = n_0 (r_s / r)$$

$$\tau_\tau(R_1, R_2) = \int_{R_1}^{R_2} n(r) \sigma_T dr = n_0 \sigma_T r_s \ln(R_2 / R_1)$$

$$L_{\text{disc}} = \begin{cases} \epsilon \dot{m} \mathcal{M} L_{\text{Edd}} & \text{for } \dot{m} \gtrsim 0.1 \\ \epsilon \dot{m}^2 \mathcal{M} L_{\text{Edd}} & \text{for } \dot{m} \lesssim 0.1 \end{cases}$$

$$U_{\text{sc}} = \frac{L_{\text{disc}} \tau_T}{4\pi R_2^2 c}$$

$$U_{\text{ext}} = \Gamma^2 U_{\text{sc}}$$

Accretion Power of the source:

$$P_{\text{acc}} = \dot{m} \mathcal{M} L_{\text{Edd}}$$

Magnetic Field

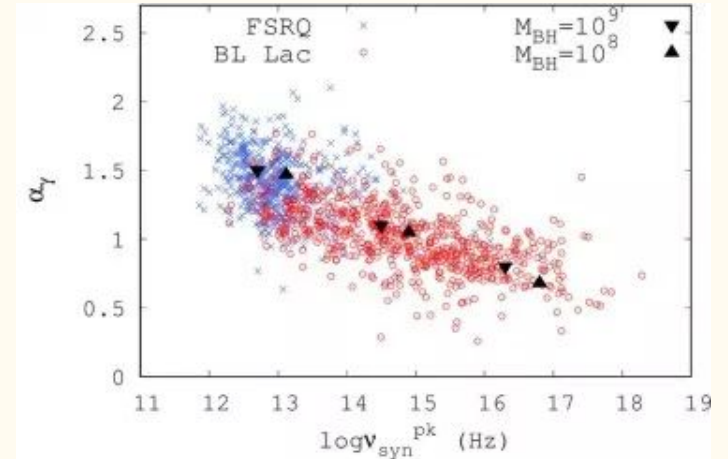
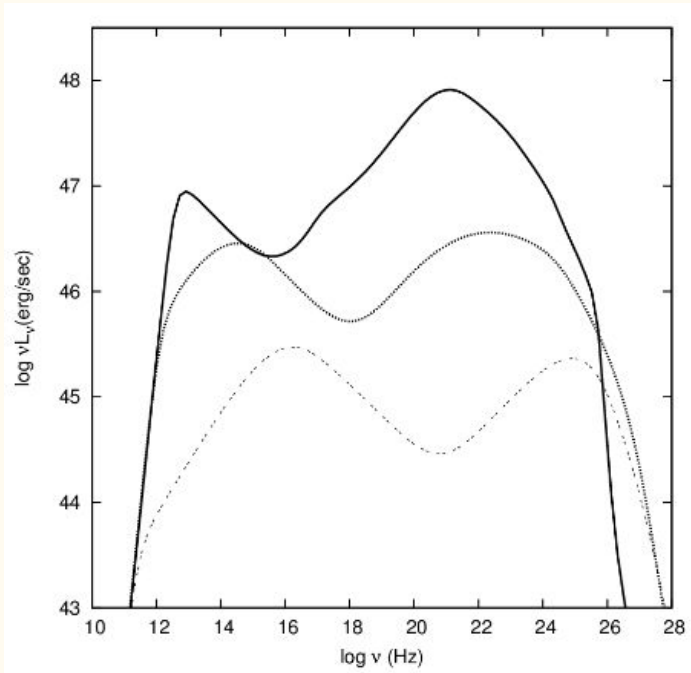
$$U_{B_0} = \frac{\eta_b P_{\text{acc}}}{4\pi(3r_s)^2 c}, \quad B = B_0 \left(\frac{z_0}{z} \right)$$

Electron Injection

$$Q_e = \begin{cases} k_{e1} \gamma^{-s} & \text{for } \gamma_{\text{min}} \leq \gamma \leq \gamma_{\text{br}}, \\ k_{e2} \gamma^{-q} e^{-\gamma/\gamma_{\text{max}}} & \text{for } \gamma_{\text{br}} \leq \gamma \leq \gamma_{\text{max}}, \end{cases}$$

$$L_{\text{inj}}^e = m_e c^2 \int_{\gamma_{\text{min}}}^{\gamma_{\text{max}}} Q_e(\gamma) \gamma d\gamma = \eta_e P_{\text{acc}}$$

$$\gamma_{\text{br}} = \frac{3m_e c^2}{4\sigma_\tau c t_{\text{dyn}} U_{\text{tot}}}$$



\dot{m}	$B(G)$	$U_{\text{ext}} \left(\frac{\text{erg}}{\text{cm}^3} \right)$	$L_e^{\text{inj}} \left(\frac{\text{erg}}{\text{sec}} \right)$	γ_{br}	Blazar Class
-0.5	-0.3	-1.4	45.2	2.3	FSRQ
-1.5	-0.8	-4.6	44.2	3.3	LBL
-2.5	-1.3	-7.6	43.2	6.5	HBL

values are in logarithmically scale

$$U_B \propto \frac{\dot{m}}{\mathcal{M}},$$

$$U_{\text{ext}} \propto U_{\text{sc}} \propto \frac{\dot{m}^{\alpha+1}}{\mathcal{M}} \quad (\alpha = 1 \text{ for } \dot{m} \geq 0.1 \text{ and } \alpha = 2 \text{ for } \dot{m} < 0.1),$$

$$\gamma_{\text{br}} \propto \dot{m}^{-1} (1 + \dot{m}^\alpha)^{-1},$$

$$L_e^{\text{inj}} \propto \dot{m} \mathcal{M}$$

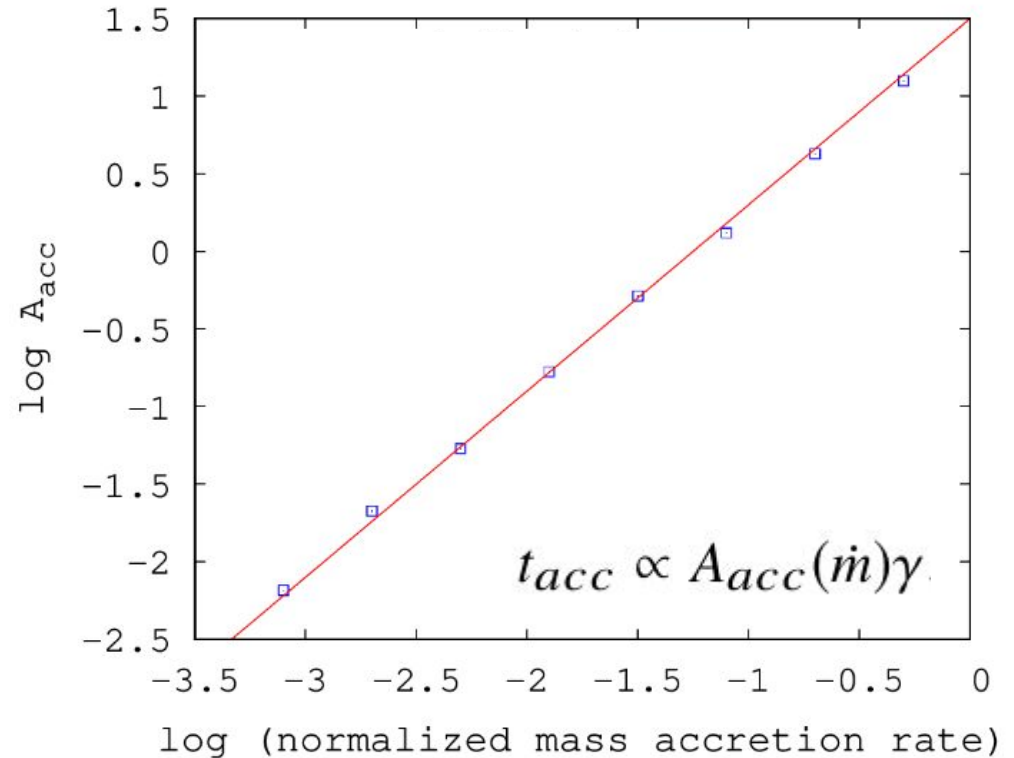
$$\nu_{\text{pk}}^{\text{syn}} \propto \mathcal{M}^{-1/2} \dot{m}^{-3/2} / (1 + \dot{m}^\alpha)^2$$

Particle Acceleration

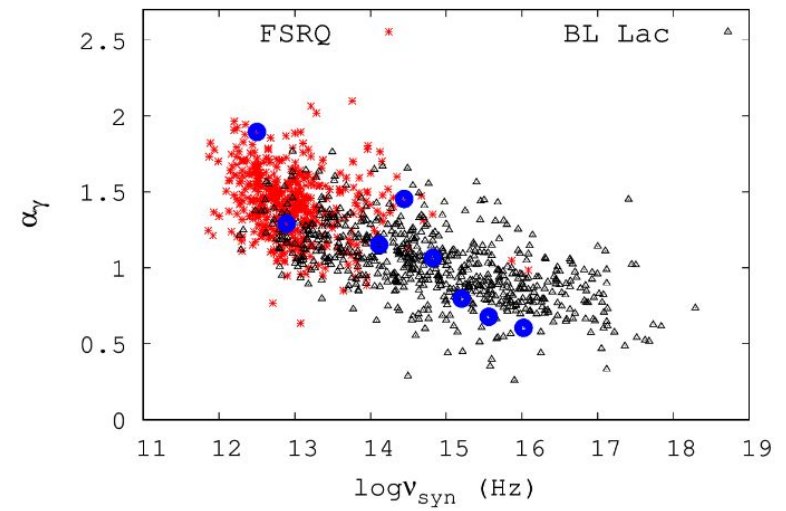
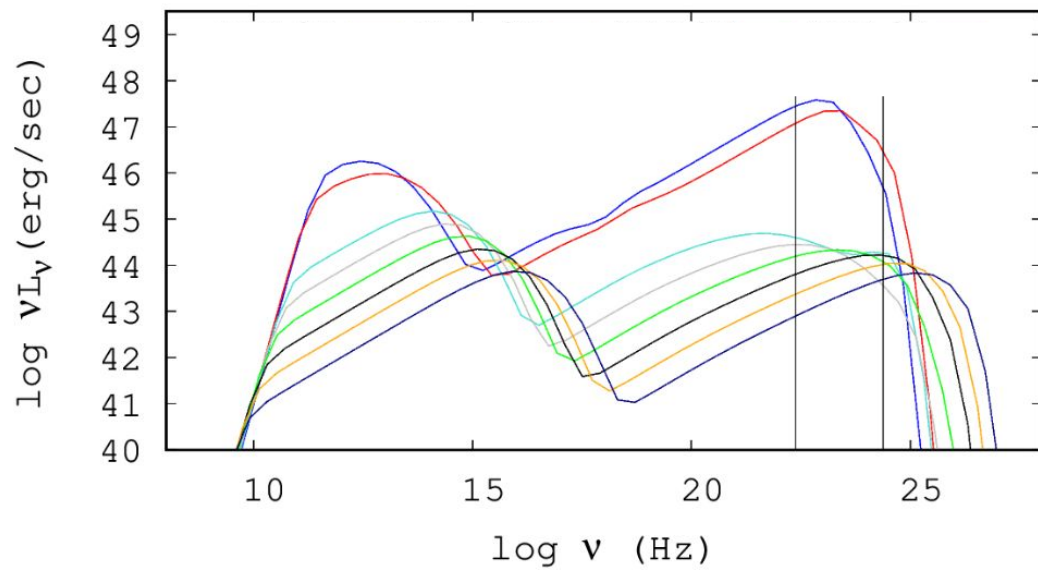
$$\frac{\partial n_e(\gamma, t)}{\partial t} + \frac{n_e(\gamma, t)}{t_{esc}(\gamma)} + \frac{\partial}{\partial \gamma} \left(\frac{\gamma}{t_{acc}(\gamma)} \right) n_e(\gamma, t) = \mathcal{L}_e(\gamma, t)$$

$$t_{accFI} \geq 6 \left(\frac{c}{u_s} \right)^2 \frac{\lambda}{c} \simeq 6 \frac{r_g c}{u_s^2},$$

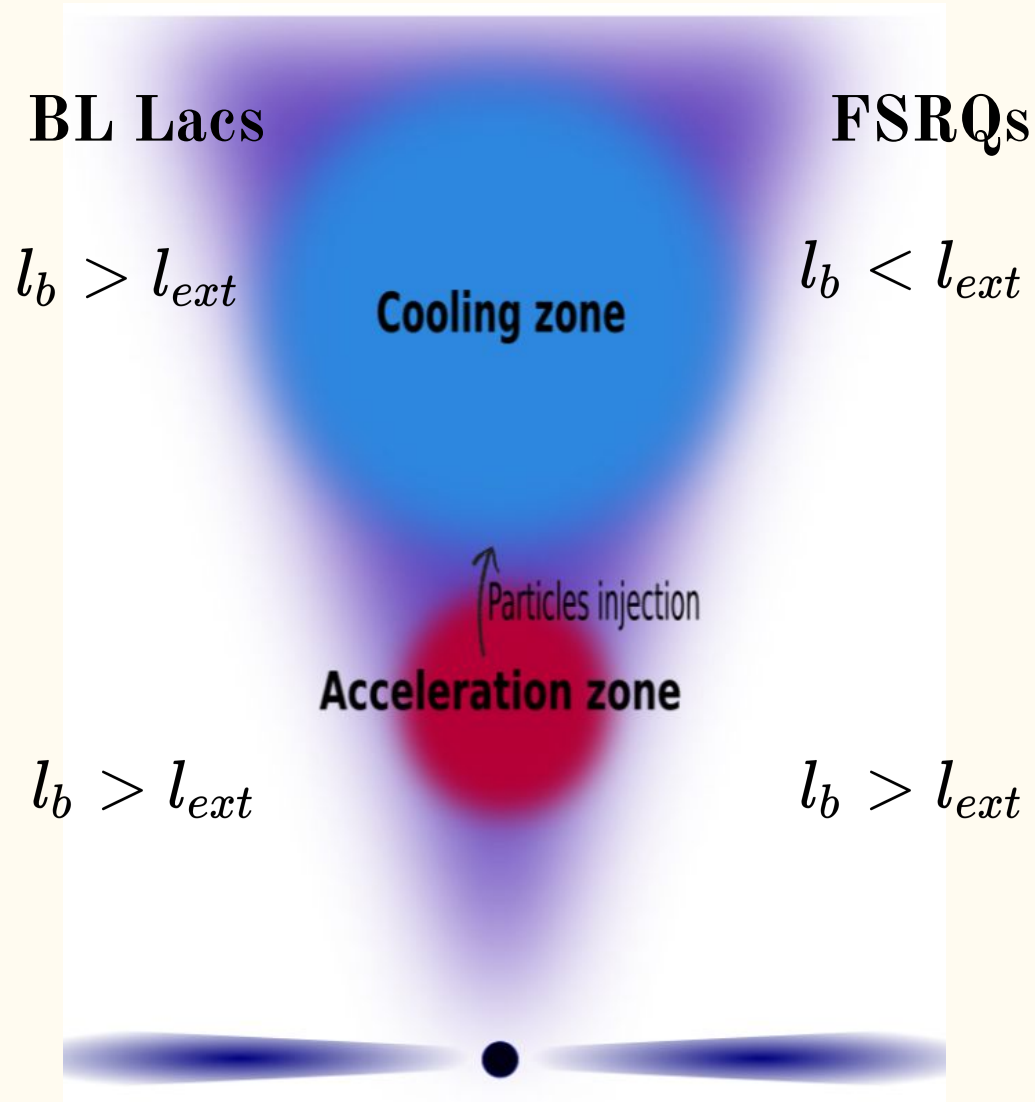
$$r_g = \frac{\gamma m c^2}{eB}.$$



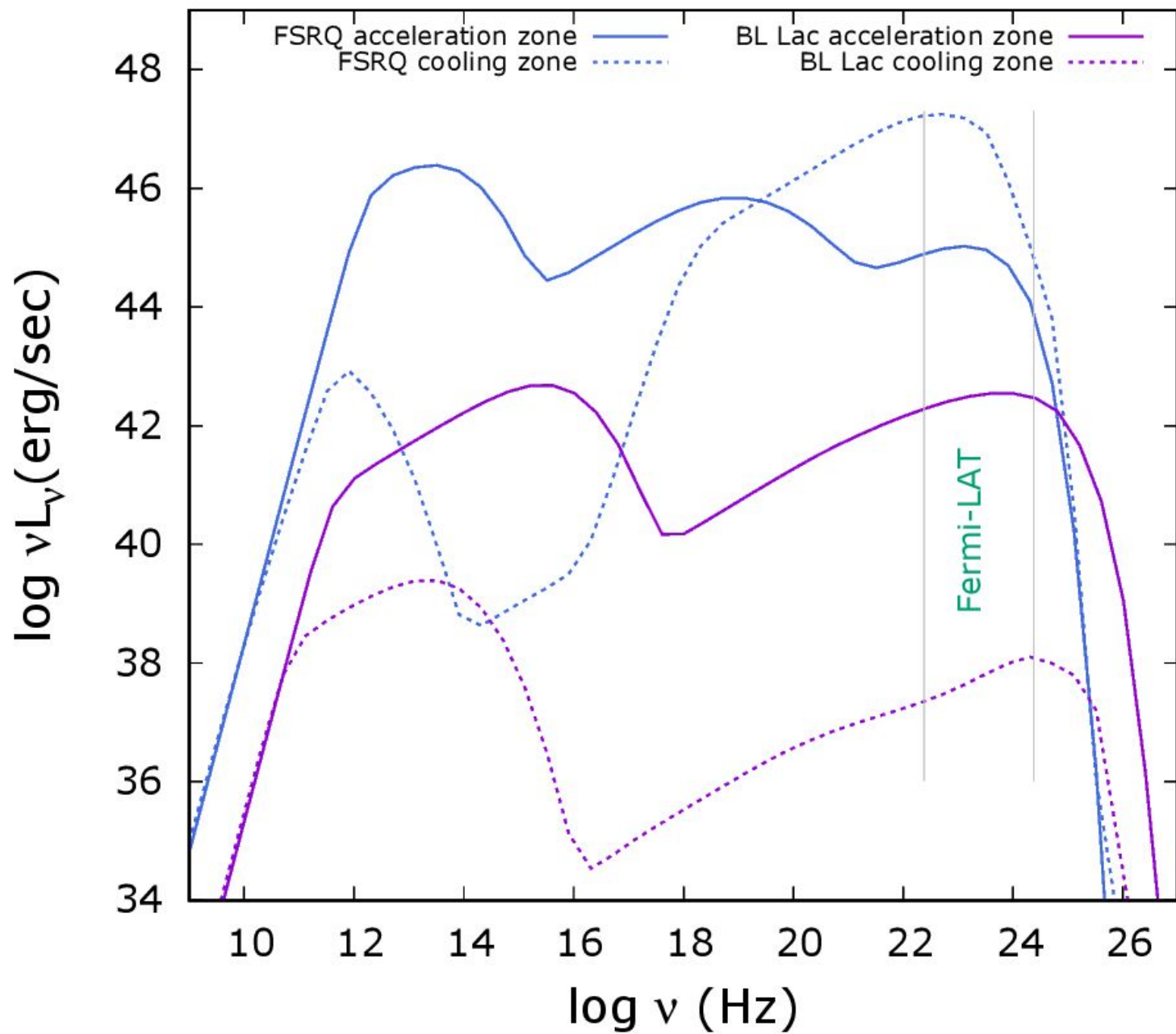
We assume next that the particles gain energy and the electrons energy distribution is calculated self-consistently

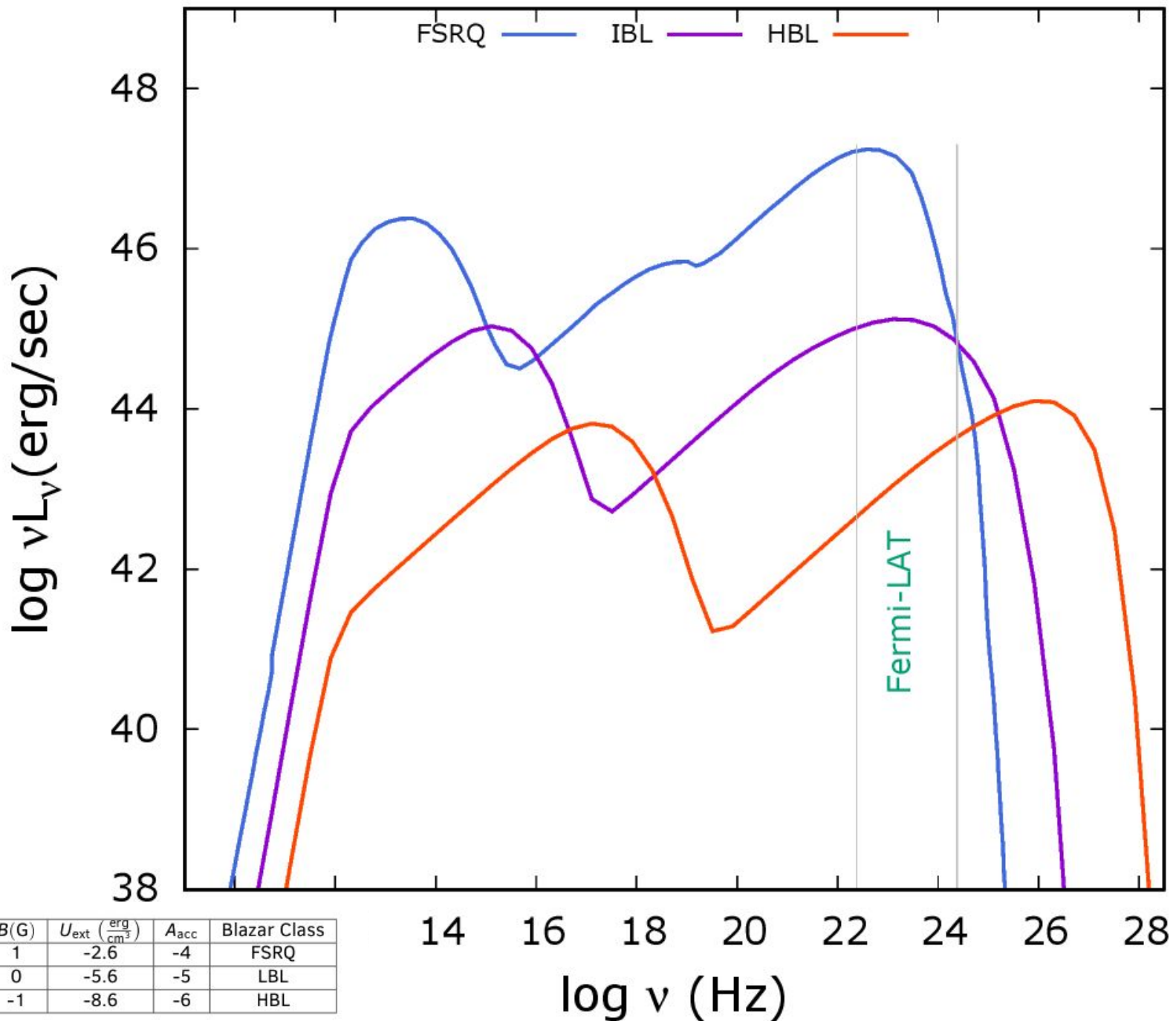


The idea:

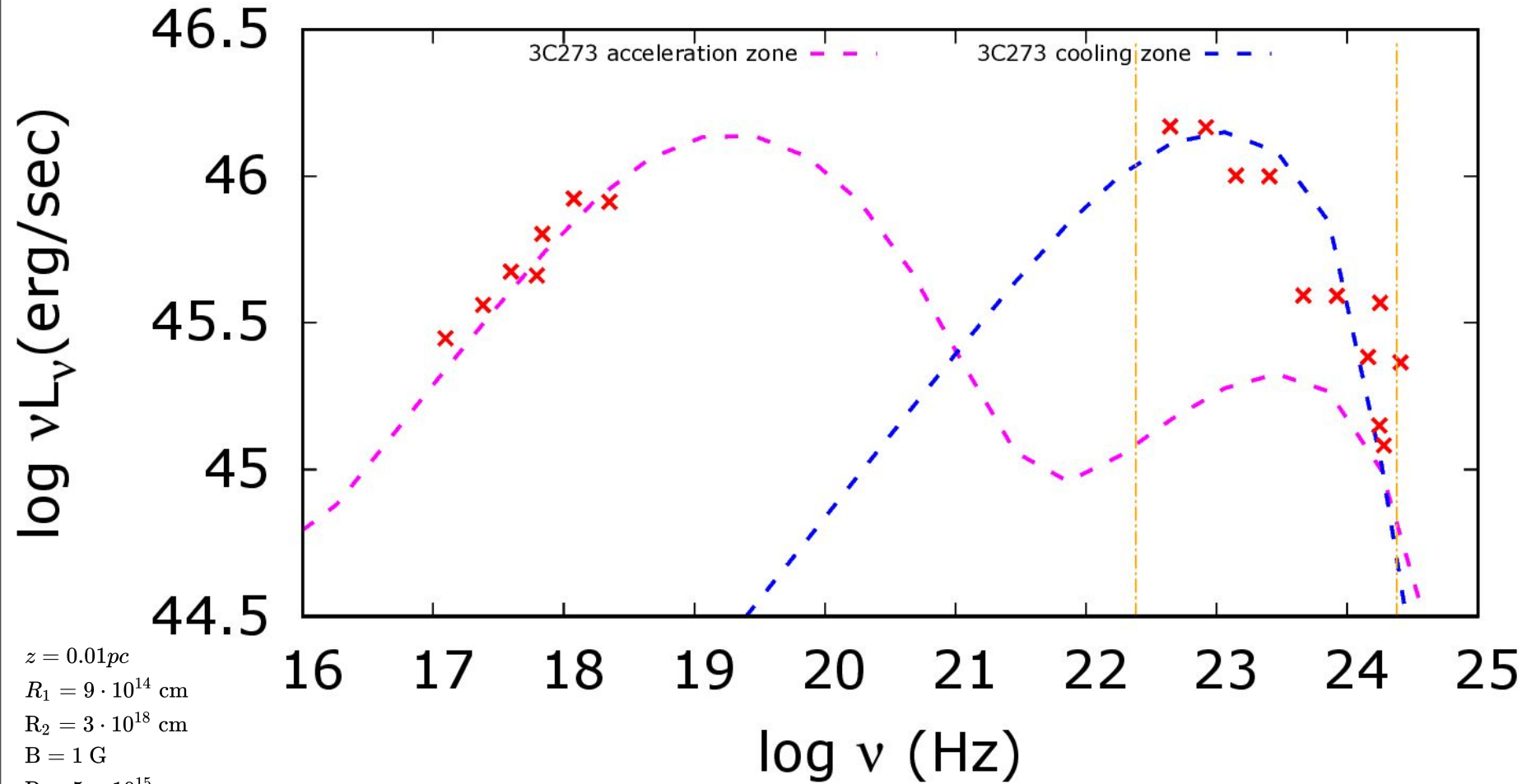


$$l_b = \frac{\sigma_\tau R_b U_b}{m_e c^2}, \quad l_{ext} = \frac{\sigma_\tau R_b U_{ext}}{m_e c^2} \quad B \propto 1/z, \quad U_{ext} = constant$$





Application to 3C273



$z = 0.01 pc$
 $R_1 = 9 \cdot 10^{14} \text{ cm}$
 $R_2 = 3 \cdot 10^{18} \text{ cm}$
 $B = 1 \text{ G}$
 $R = 5 \times 10^{15} \text{ cm}$
 $A_{\text{acc}} = 0.0001$
 $U_{\text{ext}} = 10^{-2.6} \frac{\text{erg}}{\text{sec}}$
 $\Gamma = 30$
 $\delta = 15$
 $T_{\text{disk}} = 3 \cdot 10^5 \text{ K}$

Take home messages

- MHD Accretion Disk Winds are fundamental in reproducing the LAT Blazar phenomenology (Blazar Sequence) which appears to be a one parameter family.
- We obtain the theoretical Blazar Sequence by varying only one parameter, the mass accretion rate .
- The spread of the distribution depends on the other parameters.



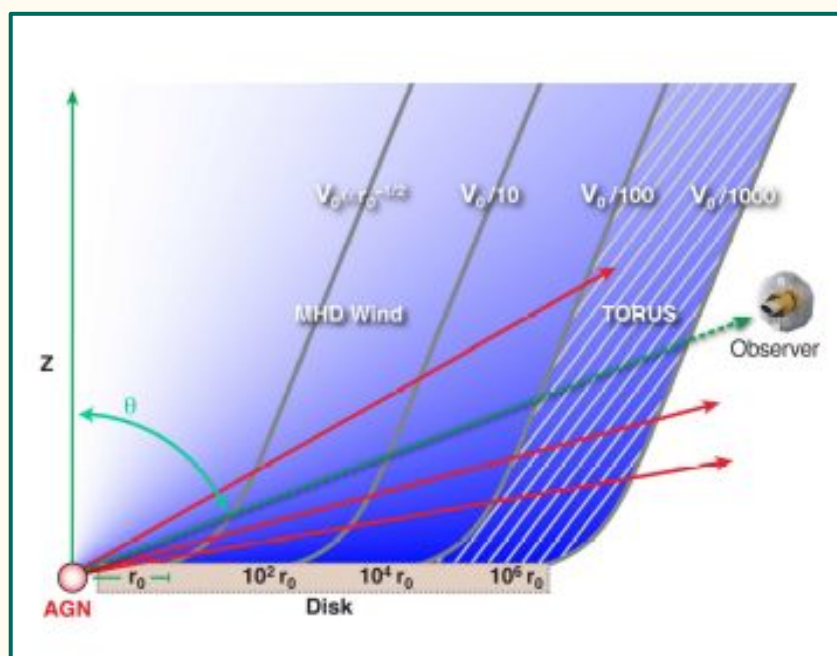
Thank you!

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Back up slides

MHD

Accretion Disk Winds

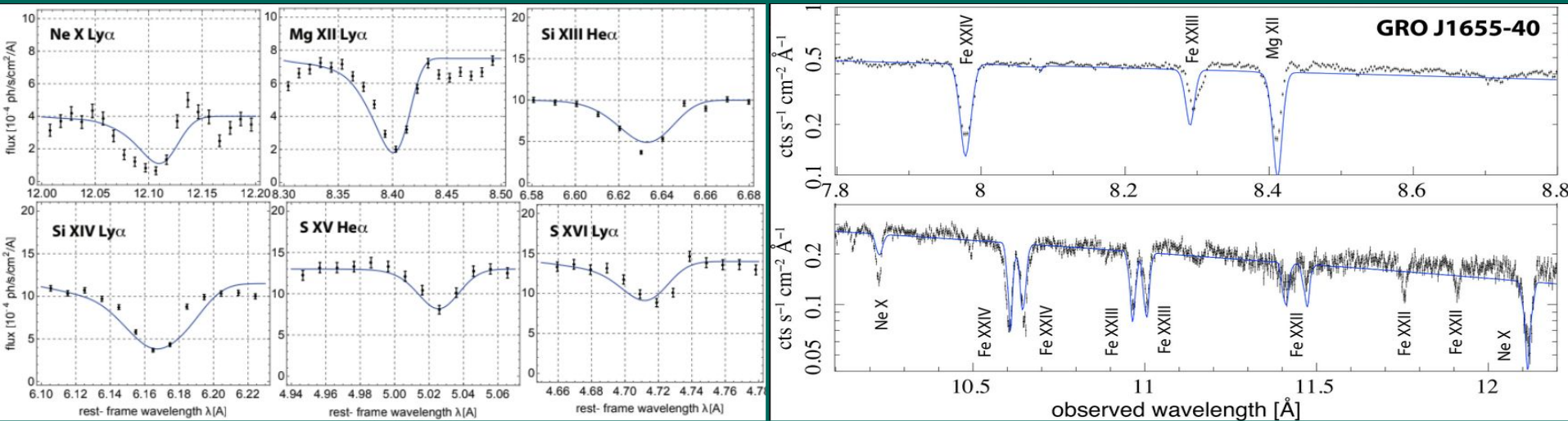


- Winds driven by an accretion disk threaded by a poloidal magnetic field.
- At latitudes above the Alfvén point the field lines become toroidal and the flow is almost radially out.
- The magnetic field permeates the entire disk, out to $\sim 10^6 R_s$

Contopoulos & Lovelace, 1994
Fukumura et al., 2010

Galactic and extragalactic applications of the wind

AGN



Fukumura et al., 2018,2019

