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Tracing the origin of low diffusivity and CR bubbles around sources

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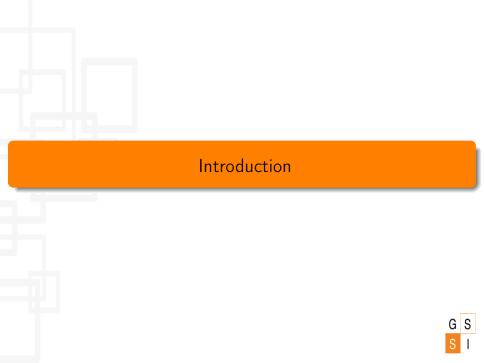
Content

Introduction

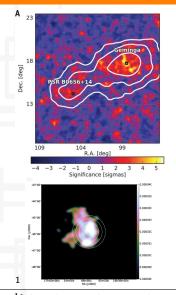
Formation of CR bubbles

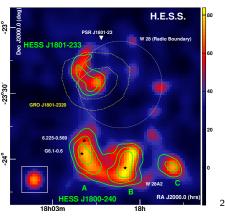
Conclusion





Motivation





 Hints for strongly reduced diffusion coefficient observed near SNR, stellar clusters and pulsars



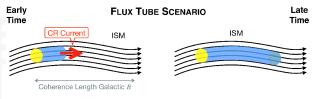
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¹[Abeysekara et al., 2017; Aharonian et al., 2019]

²[Aharonian et al., 2008]

Standard picture of CR escape

- interstellar magnetic field is coherent on scales of ~ 10 -50 pc
- particles mainly diffuse parallel to the magnetic field



⇒ on these scales particles fill a flux tube and the problem can be regarded 1D

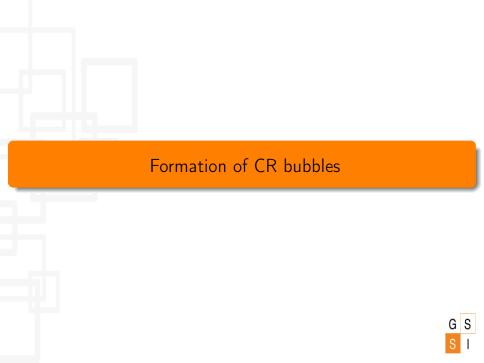
- flux given by: $\phi_{CR}(E > E_0) = n_{CR}(E > E_0)v_D = \frac{L_{CR}}{2\pi R^2 \Lambda E_0}$
- In this configuration CRs will excite a resonant streaming instability, hindering their escape

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Less efficient for high-energy particles ³



³[Nava et al., 2016, D'Angelo et al. 2016]



Bell instability ⁴

Did we miss something?





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Bell instability 4

Did we miss something? ⇒ non-resonant streaming instability

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- large growth rate γ_{max}
- on scale $k^{-1} \ll r_L \Rightarrow$ does not affect the CR current at first
- until saturation at $\sim 5-10\gamma_{max}^{-1}$, then cascades to larger scales





Bell instability 4

Did we miss something? ⇒ non-resonant streaming instability

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$$D = \frac{1}{3} \frac{vr_L}{P(k_{\rm res})} \sim \frac{B_0^2}{\delta B^2}$$
 \rightarrow affects transport and enables strong particle scattering

Condition:

$$\frac{\phi_{CR}(E>E_0)}{c}E_0\gg \frac{B_0^2}{4\pi}$$

This instability is often excited at SNR shocks





Source in the ISM

Bell condition for SNR:

$$rac{4\pi\phi_{CR}}{cB_0^2}pprox 100$$
 ,

,
$$\gamma_{\it max}^{-1} pprox 1.1 ({
m E}/2.5 {
m TeV})$$
 yr

Easy to see: Flux at SNR shock = flux of escaping particles, the condition holds at the shock \Rightarrow it holds in the flux tube



Source in the ISM

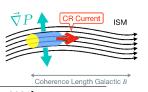
Bell condition for SNR:

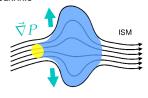
$$\frac{4\pi\phi_{\mathit{CR}}}{cB_0^2}\approx 100$$

$$\gamma_{\it max}^{-1} pprox 1.1 (E/2.5 {
m TeV}) \, {
m yr}$$

- Easy to see: Flux at SNR shock = flux of escaping particles, the condition holds at the shock \Rightarrow it holds in the flux tube
- mean free path $\lambda = \frac{3D}{V} \approx 1 \cdot E_{GAV}^{1/2} \, \text{pc} \Rightarrow \text{high-energy CR escape ballistically}$
- Flux $\phi_{CR}(E > E_0) = n_{CR}(E > E_0)v_D$ conserved, but v_D decreases as particles start to transition to diffusive behavior
- $\bullet \Rightarrow P_{CR} \gg P_{gas} \rightarrow$ tube will expand in transverse direction, breaks 1D geometry

BUBBLE SCENARIO

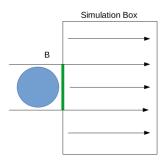




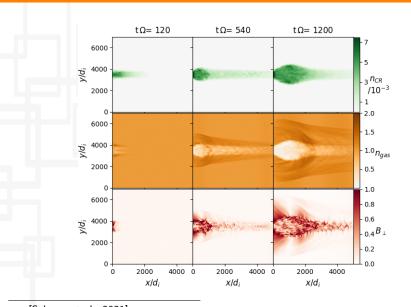


Simulation

- Hybrid particle-in-cell simulation with dHybridR
- Solve Maxwell equations and equations of motion for macroparticles
- Electromagnetic fields due to moving particles

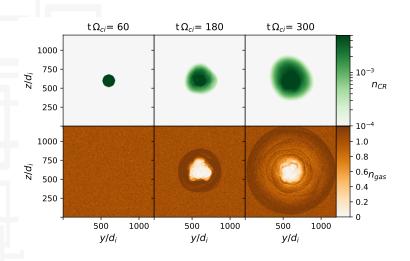


2D Simulation





3D Simulation







Conclusion

new insights about escape of CRs from their sources:

- current of escaping particles generate a non-resonant instability which slows down their escape
- leads to formation of CR bubbles around sources with reduced diffusivity
- Important implications:
 - enhanced γ -ray emission from circumsource region
 - accumulated grammage of trapped CRs might be significant for secondary-to-primary ratios
 - electrons get trapped as well so that energy losses become important

