

# Probing cosmic ray distribution around Cygnus OB2

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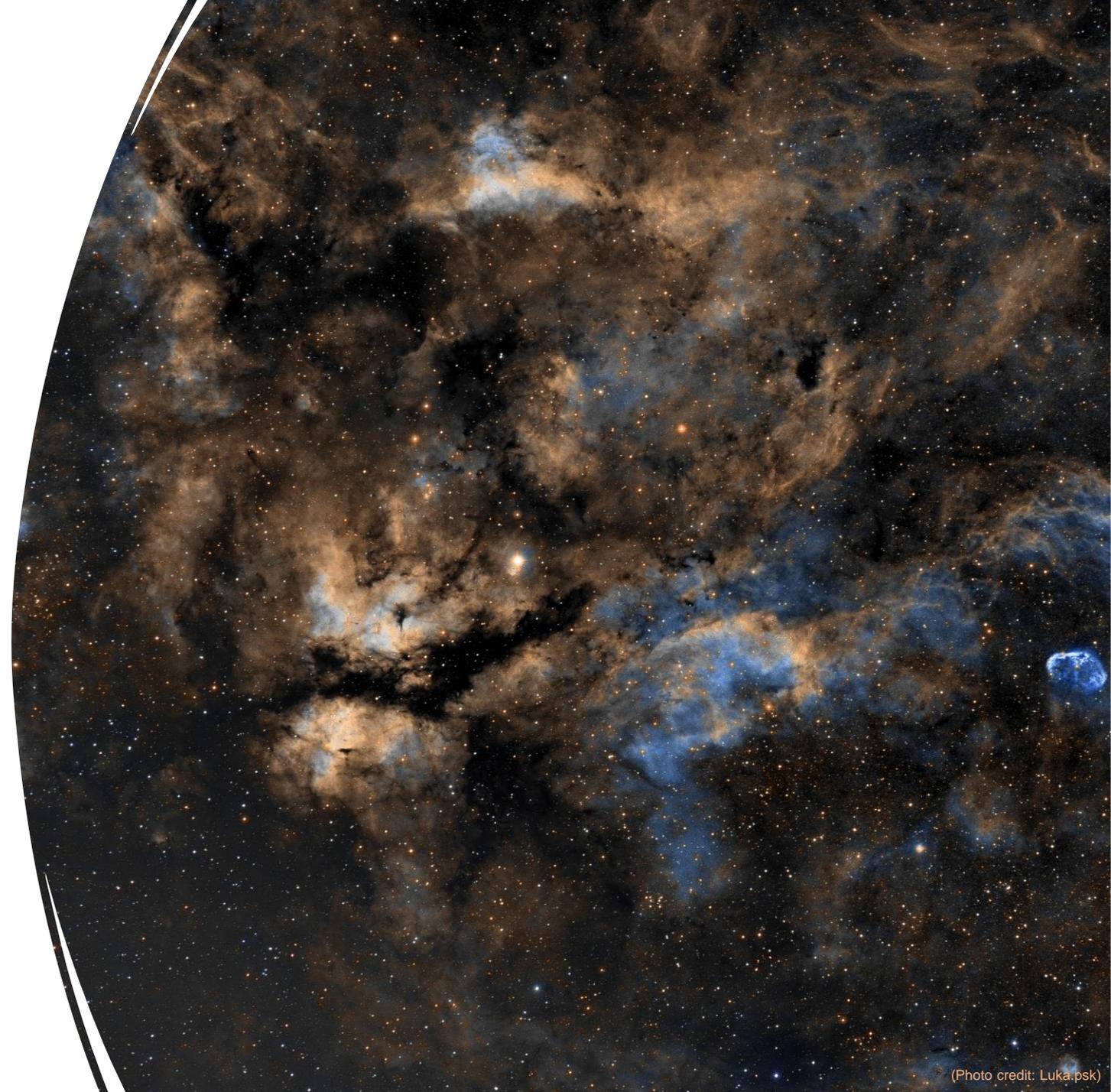
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## Outline

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- Cygnus OB2
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  - Modified solution
- Interstellar medium around OB2
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# Introduction

**Massive star clusters** (MSCs) are thought to be a possible class of cosmic ray (CR) accelerators powered by the strong winds blown by the stars inside the cluster.

High energy and very-high energy  $\gamma$ -ray emission has been observed in the direction of several MSCs, such as: Cygnus OB2, Westerlund 1, Westerlund 2, NGC3603, ...

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Cygnus OB2



Westerlund 1

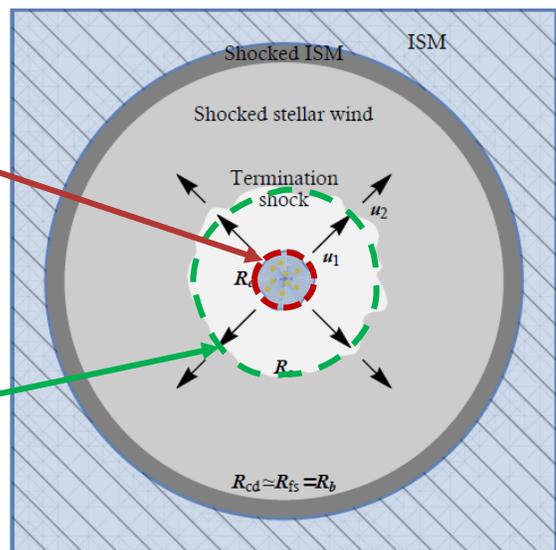


The acceleration mechanism is still under debate.

For example, two possible cases are:

Wind-wind acceleration  
inside the cluster core

Acceleration at the  
MSC's wind termination  
shock



Diverse models will produce different distributions of CRs around the cluster. Depending on the distribution of the interstellar medium (ISM) in the neighborhood of the cluster, the  $\gamma$ -ray morphology and spectrum may vary

## Objective:

Consider the case of **Cygnus OB2** and compare with available data the expected  $\gamma$ -ray emission (**spectral energy distribution and spatial morphology**) assuming the model where CRs are accelerated at the cluster wind's termination shock

# Cygnus OB2

**Cygnus OB2** is one of the most massive MSC in the Milky Way, hosting  $\approx 170$  OB stars (Wright et al. 2015).

OB2 is located in the Cygnus X star forming complex, positioned tangent to the local spiral arm ( $l=80.22^\circ; b=0.77^\circ$ )

## Cygnus OB2 parameters

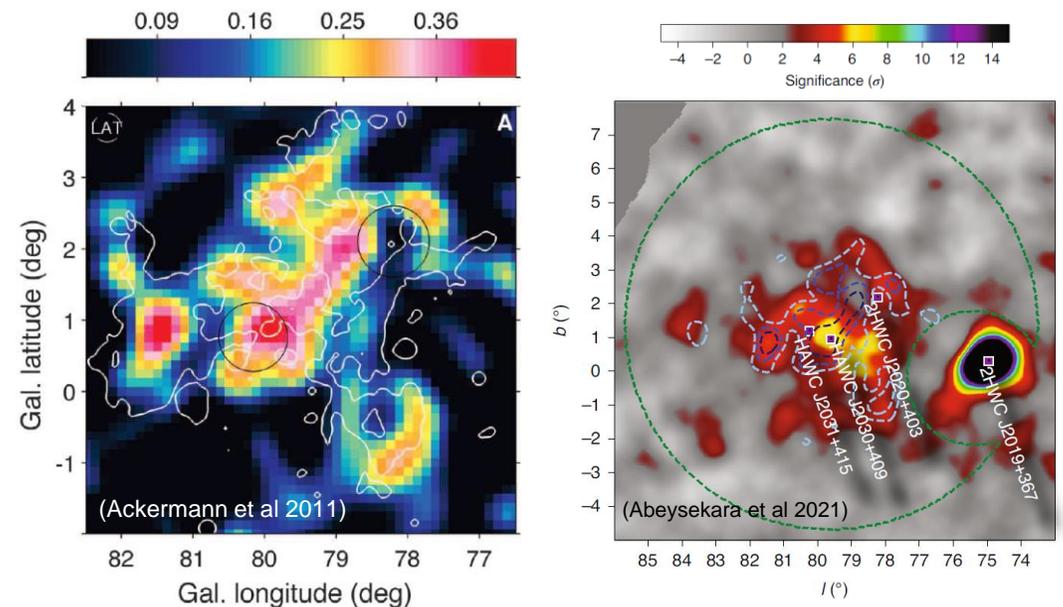
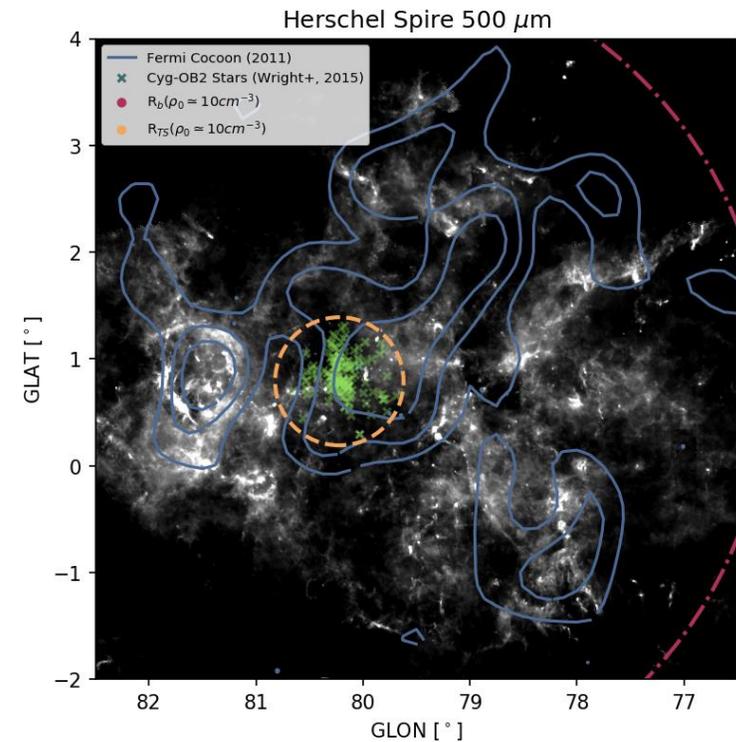
$$L_{\text{wind}} = 2 \times 10^{38} \text{ erg/s}; \quad dM/dt = 10^{-4} M_{\text{sun}}/\text{yr}; \quad d_{\text{OB2}} = 1.4 \text{ kpc}$$

$$u_1 \approx 2500 \text{ km/s}; \quad \rho_H = 10/\text{cm}^3; \quad t_{\text{age}} = 3 \text{ Myr}$$

$$R_{TS} = 0.7 \cdot L_w^{-1/5} \dot{M}^{1/2} u_1^{1/2} \rho_H^{-3/10} t_{\text{age}}^{2/5} \simeq 16 \text{ pc}$$

$$R_b = 0.76 \cdot \left( \frac{L_w}{\rho_H} \right)^{1/5} t_{\text{age}}^{3/5} \simeq 98 \text{ pc}$$

Extended  $\gamma$ -ray emission has been detected by several experiments:  
Fermi (2011), Argo (2014),  
HAWC (2021), LHAASO (2021)



# CRs distribution function

## Steady state solution

We use the model developed by [Morlino et al. \(2021\)](#) of **CR accelerated at the winds' termination shock** from MSC to obtain the CR distribution function around OB2.

This model considers a steady state solution for CRs injected at the termination that escape from the system considering both advection and diffusion.



$$f(r < R_{TS}; E) = f_{TS}(E) \cdot \exp\left[-\frac{u_1(R_{TS} - r)}{D_1}\right]$$

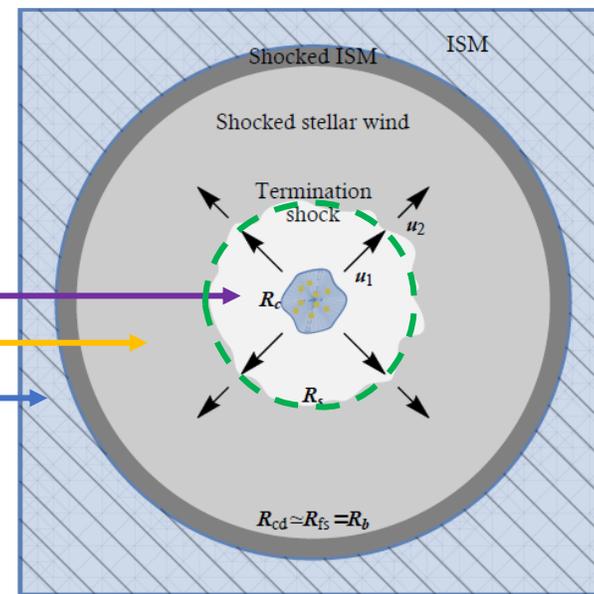
$$f(R_{TS} < r < R_b; E) = f_{TS}(E)\Gamma_1 + f_{gal}(E)\Gamma_2$$

$$f(r > R_b; E) = f(R_b; E) \frac{R_b}{r} + f_{gal}(E) \left(1 - \frac{R_b}{r}\right)$$

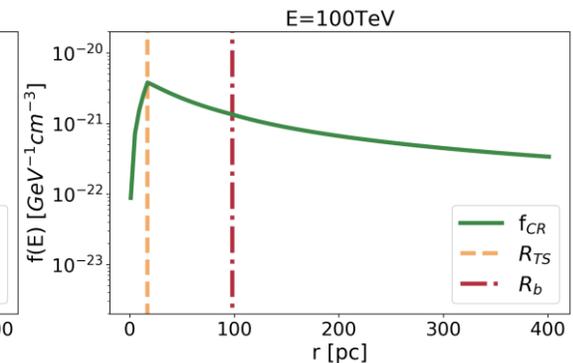
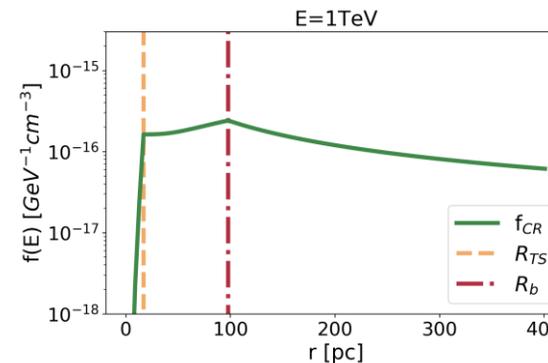
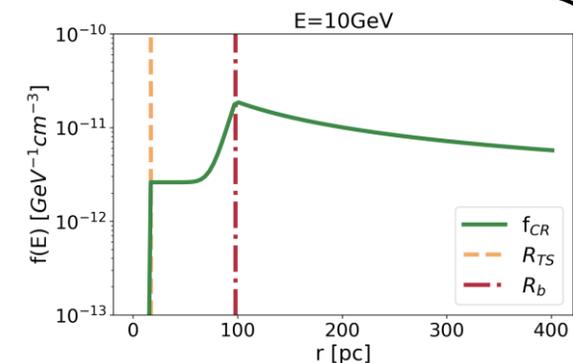
Where  $\Gamma_1$  and  $\Gamma_2$  are function depending on  $D_2$ ,  $D_{ism}$ ,  $u_2$ ,  $R_{TS}$  and  $R_b$   
The distribution function at the **termination shock** is modeled as:

$$f_{TS}(E) = A \left(\frac{E}{1 \text{ TeV}}\right)^{-a} \exp\left[-\left(\frac{E}{E_{max}}\right)^b\right]$$

The parameters of  $f_{TS}$  are chosen to best reproduce  $\gamma$ -ray observations

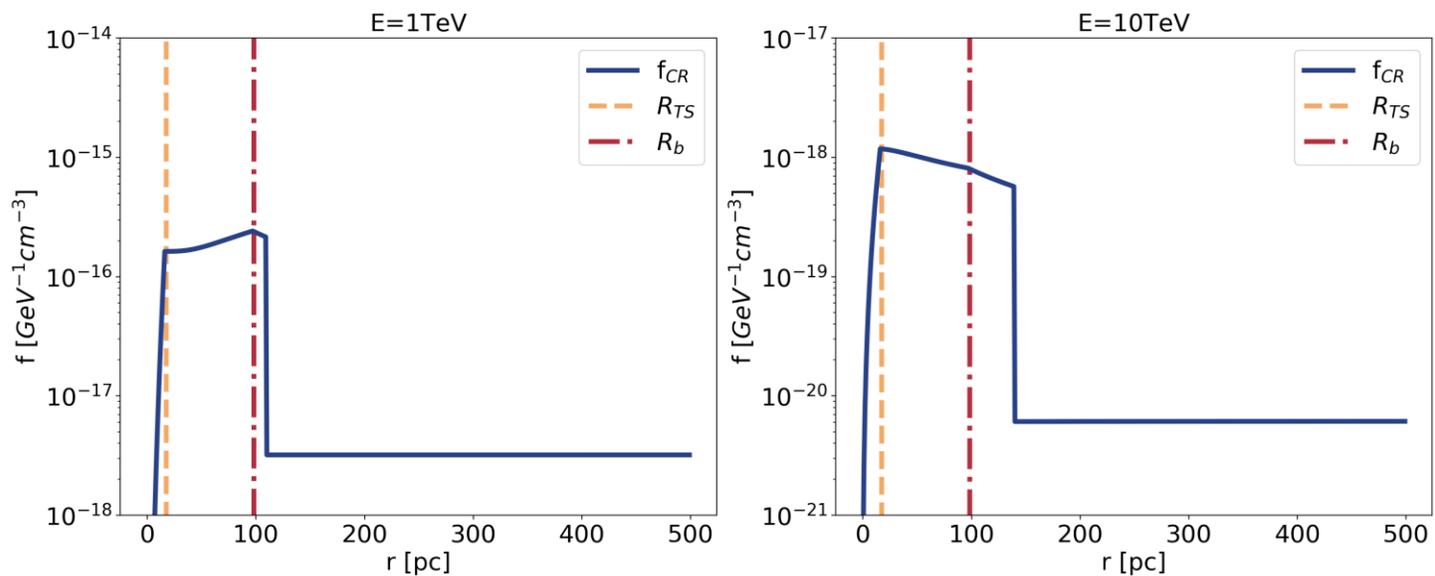
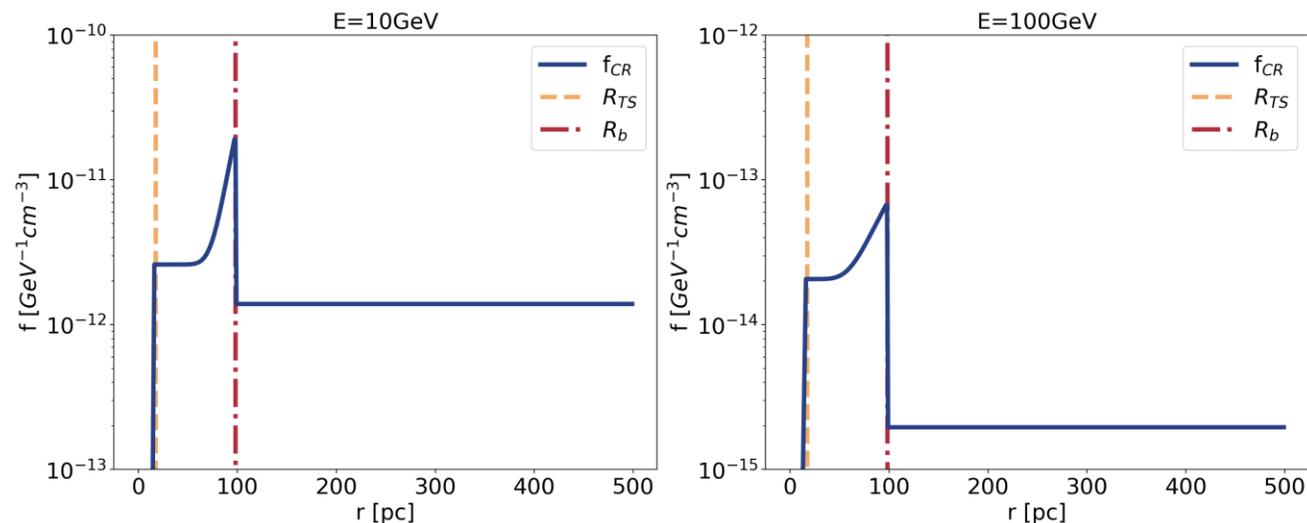


Steady state radial distribution of CRs accelerated by Cygnus



# CRs distribution function

## Modified solution

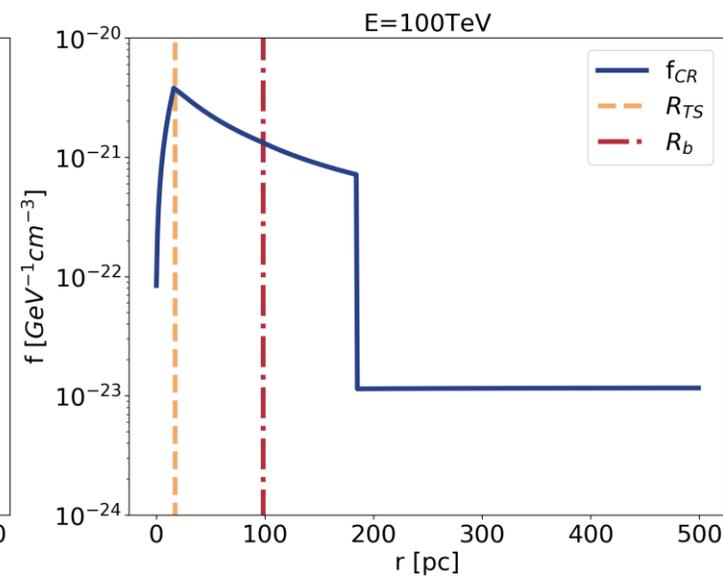


**In the steady state model CRs diffuse up to infinity.**  
 $f(r,E)$  must be truncated at a given  $R_{\text{cut}}(E)$  so that the total injected energy by OB2 is conserved.

Assuming no energy losses,  $R_{\text{cut}}(E)$  can be calculated by requiring:

$$\int_0^{R_{\text{cut}}(E)} 4\pi r^2 [f(r, E) - f_{\text{gal}}(E) \Theta(r - R_{\text{TS}})] dr = 4\pi R_{\text{TS}}^2 t_{\text{age}} u_2 f_{\text{TS}}(E)$$

**NB:**  $R_{\text{cut}}(E) \approx R_b$  for energies below 1 TeV.  
Low energy CRs are dominated by advection and remain trapped inside the expanding bubble



# Interstellar medium around OB2

We model the ISM around OB2 as a combination of **molecular** ( $H_2$ ) and **neutral** (HI) hydrogen.

Kinematic cuts:  $-20\text{km/s} < v < 20\text{km/s}$

**$H_2$ :**  $^{12}\text{CO}(J=1-0)$  CfA (Dame et al, 2001) Lowres.  
(using  $X_{\text{co}}=1.68 \times 10^{20} \text{ mol. cm}^{-2} \text{ K}^{-1} \text{ km}^{-1}$ )

$^{13}\text{CO}(J=1-0)$  NRO (Takekoshi et al, 2019) Highres.  
(assuming LTE +  $[H_2]/[^{13}\text{CO}]=69 \times 10^4$ )

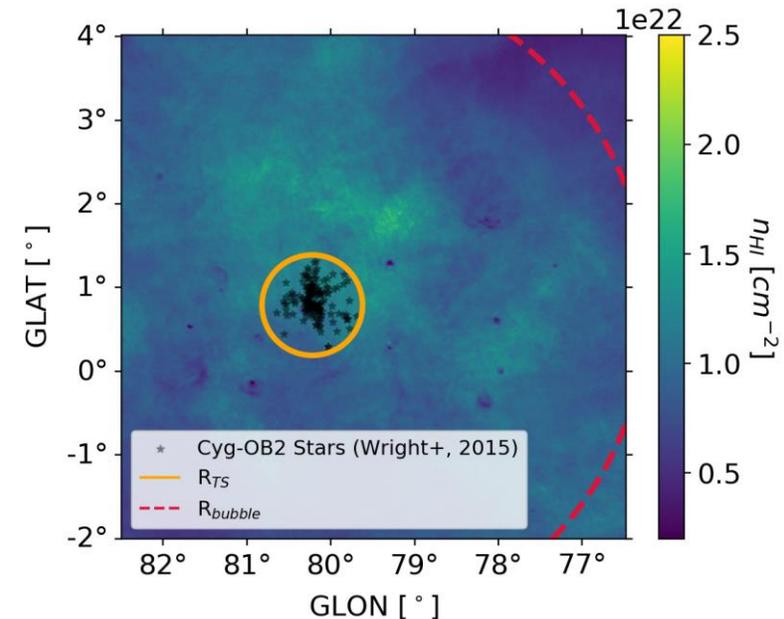
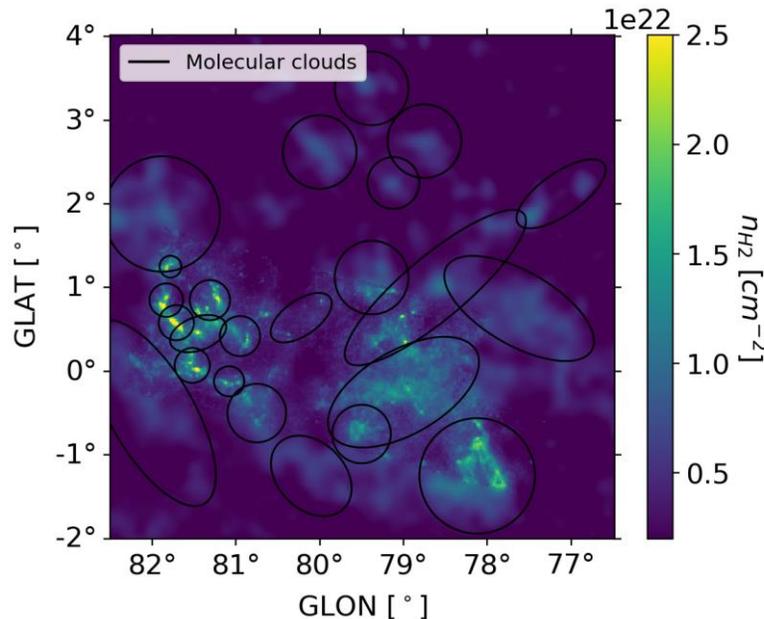
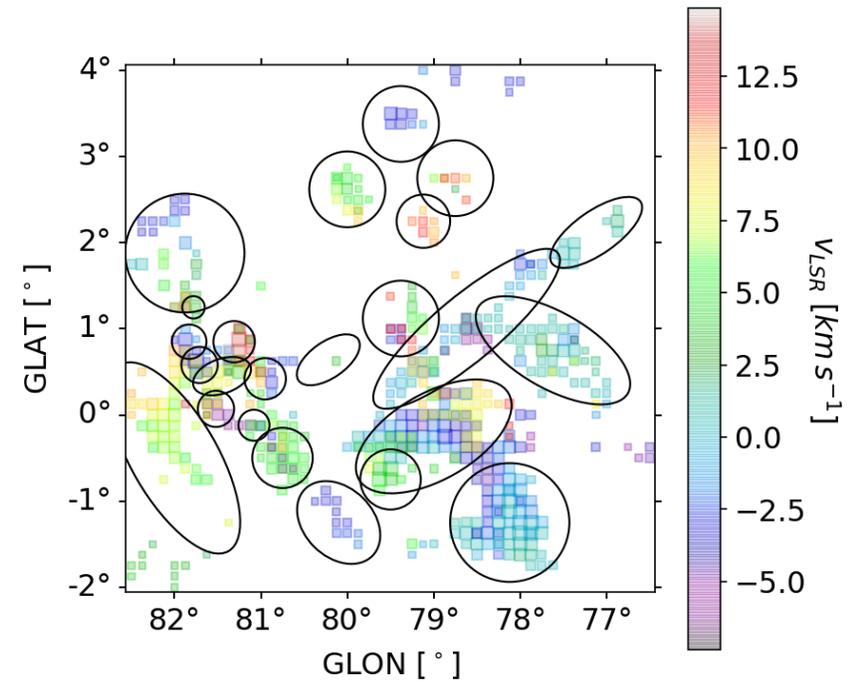
**HI:** 21cm from CGPS (Taylor et al, 2003)  
(using  $T_{\text{spin}}=150^\circ\text{K}$ )

## Two different ISM distributions are considered:

- 1) HI and  $H_2$  uniformly distributed along the line of sight in  $\pm 400\text{pc}$
- 2) Complex ISM distribution with MCs positioned at random distance from OB2 (in  $\pm 400\text{pc}$ ), and HI uniformly distributed along the line of sight for  $r > R_b$  (in  $\pm 400\text{pc}$ ). Inside  $R_b$  we consider a particle density of  $10^{-3} \text{ cm}^{-3}$

We consider a total of 24 MCs

MCs in the area are selected considering spatial shape and  $^{12}\text{CO}$  line velocity



# Results

## Spectral energy distribution

When considering the case of complex ISM distribution, we calculate the expected flux using 100 different template realizations, randomly varying MCs' position

The spectral energy distribution (SED) is extracted from a 2.2° region centered on OB2. Data points from experiments are scaled to account only the flux coming from a region of this size.

**The obtained SED fairly well fit the data point.**  
The distribution of MCs induces a little variance on the SED (factor ≈2) at ≈ 100 GeV.

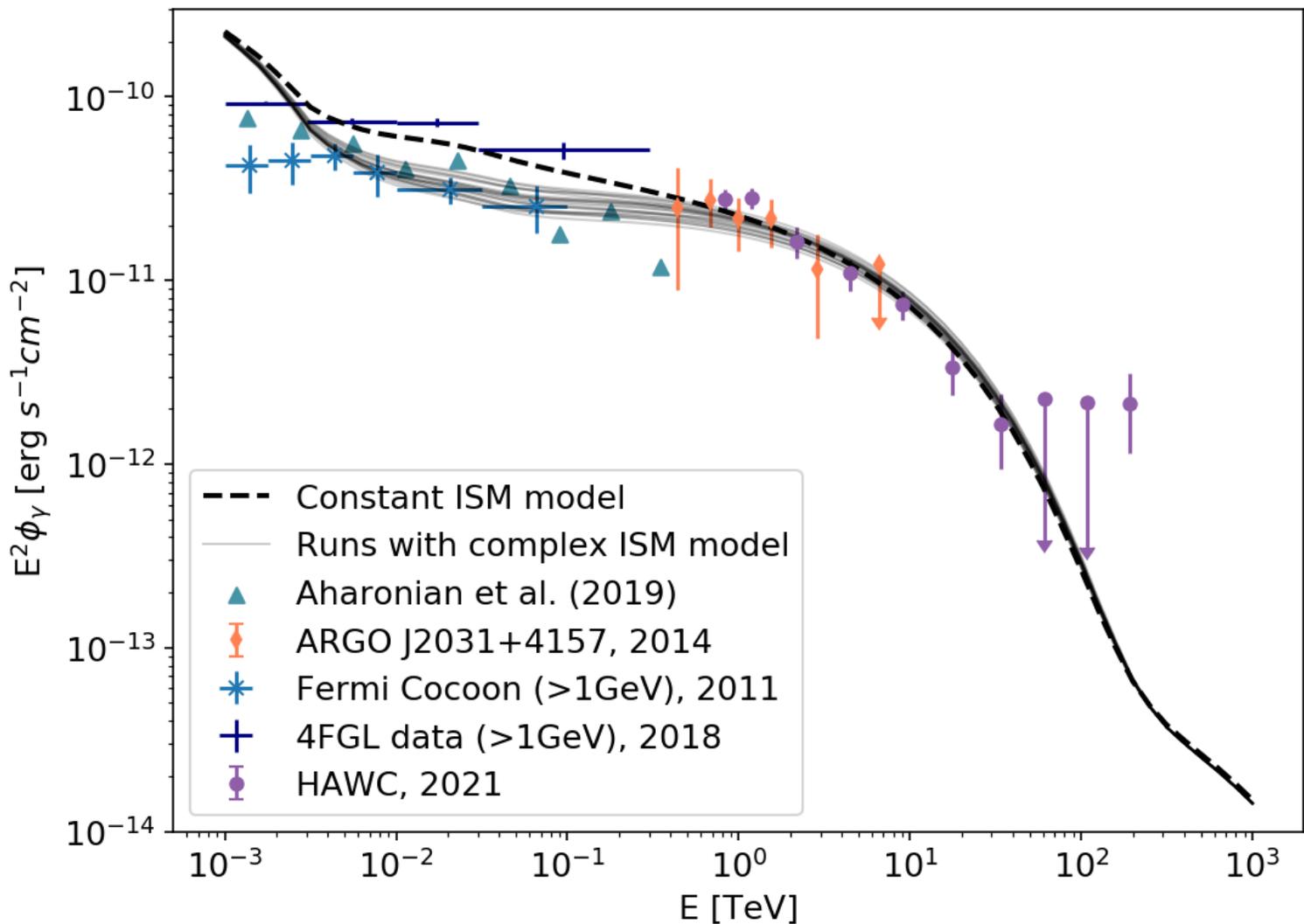
The following parameters for  $f_{\text{TS}}$  are used when computing the SEDs:

$a=2.1$ ,  $b=1$ ,  $E_{\text{max}}=100\text{TeV}$ ,  
 $A=1.64 \times 10^{-16} \text{ cm}^{-3} \text{ GeV}^{-1}$  (const. ISM case),  
 $A=0.6 \times 10^{-16} \text{ cm}^{-3} \text{ GeV}^{-1}$  (complex ISM runs)

The  $\gamma$ -ray flux from hadronic interaction is calculated using:

$$\phi_{\gamma}(l, b; E_{\gamma}) = \iint \frac{c \Omega r^2}{4\pi r^2} n_{\text{gas}}(l, b, z) f_{\text{CR}}(l, b, z; E_p) \frac{d\sigma(E_p, E_{\gamma})}{dE_p} dE_p dz$$

where  $\sigma$  is the cross section for gamma-ray production (Kafexhiu et al 2014),  $\Omega$  is the pixel size of  $(0.02^{\circ})^2$  and  $f_{\text{CR}}$  is the distribution of CRs around OB2.



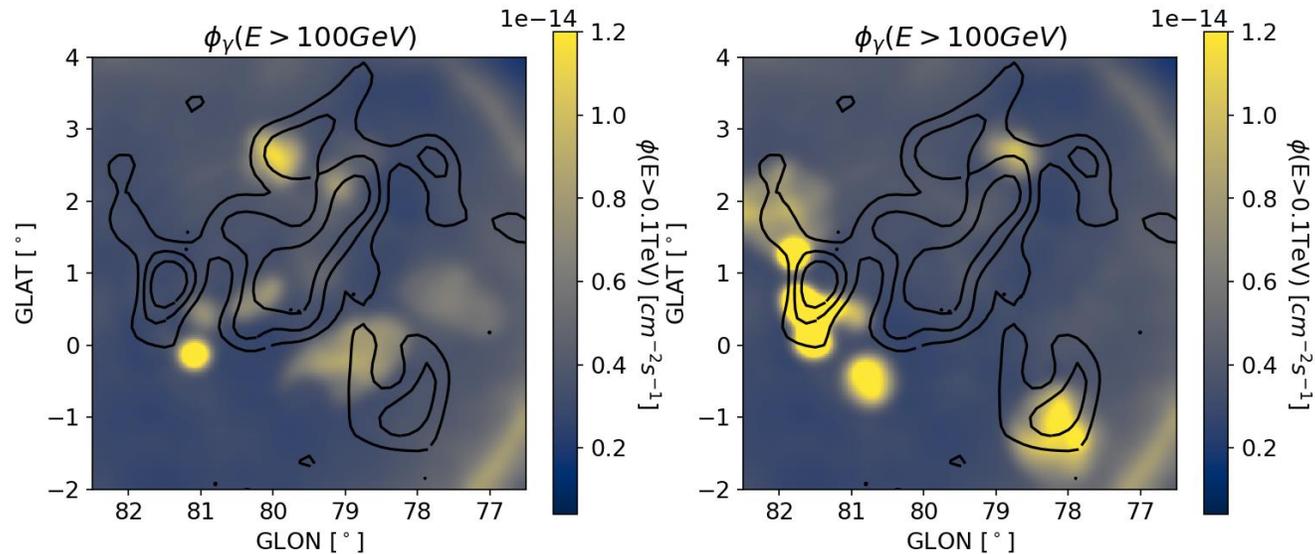
# Results

## $\gamma$ -ray spatial morphology

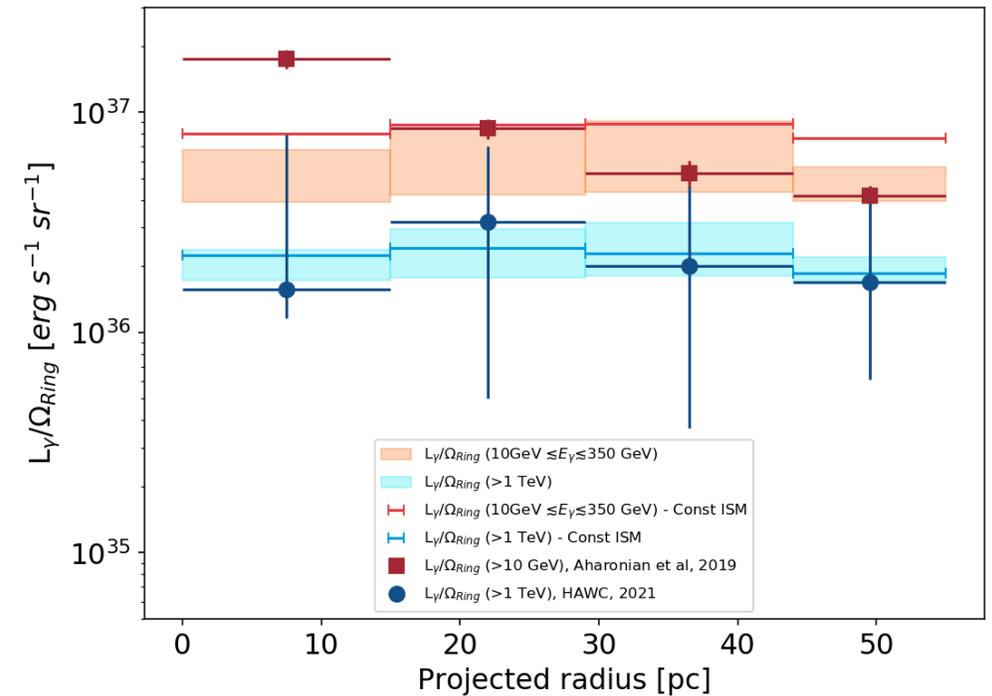
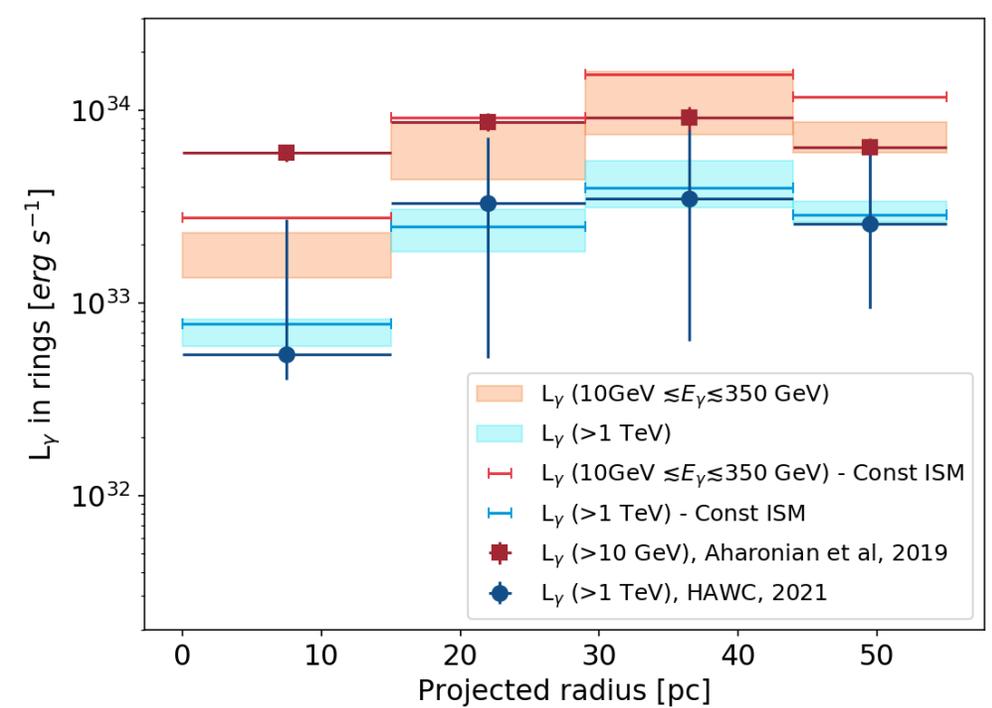
MCs position significantly impact on the observed morphology.

Radial profile extracted from 4 rings  $[0-0.6^\circ, 0.6-1.2^\circ, 1.2-1.8^\circ$  and  $1.8-2.2^\circ]$  centered on OB2 is characterized by a **flat trend**.

Radial profile in agreement with all data except for the luminosity in the first ring observed by Fermi (Aharonian et al, 2019)



Flux maps convolved with a  $0.1^\circ$  gaussian kernel



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# Conclusions

## What has been done:

- We simulated hadronic  $\gamma$ -ray emission from Cygnus OB2 assuming a model of CRs accelerated at the cluster wind's termination shock.
- We consider two possible cases of ISM distribution around the cluster, one with uniformly distributed gas, the other with MCs scattered around OB2.
- The SED adequately fit the data from Fermi-LAT to the very-high energy band observed by HAWC
- The  $\gamma$ -ray morphology at  $\approx 100\text{GeV}$  depends on the relative position between the MCs and OB2.
- The expected radial shape shows a flat trend, in agreement with HAWC data but not completely with Fermi observations.

## Possible future steps

- Consider the case where acceleration takes place inside the cluster
- Analyze the contribution from leptonic emission

