



Study of the gamma-ray state changes of PSR J2021+4026

² with Fermi-LAT

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The peculiar radio-quiet PSR J2021+4026, in the Gamma Cygni supernova remnant, is one of the brightest of the over 250 pulsars detected by *Fermi*-LAT. It is unique in being the only known isolated gamma-ray pulsar to undergo abrupt flux changes, which are also simultaneous with spin-down variations. The first change was observed by *Fermi*-LAT in October 2011, and it was followed by a recovery over a timescale of 100 days around December 2014. A second change occurred in February 2018. In the last few years, PSR J2021+4026 has been widely studied

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occurred in February 2018. In the last few years, PSR J2021+4026 has been widely studied at different wavelengths. We report our latest results on this source, based on a *Fermi*-LAT analysis of its gamma-ray variability. In particular, we have studied the changes in the spectral and timing parameters on different timescales. Our results are essential to relate the observed events to changes in the geometry of the particle acceleration regions in the pulsar magnetosphere. Therefore, this study will allow us to enhance our knowledge of this source and its behavior.

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14 **1. Introduction**

¹⁵ The young, radio-quiet γ -ray pulsar PSR J2021+4026 was discovered with the *Fermi* Large ¹⁶ Area Telescope (LAT) [1] within the shell of the Gamma Cygni supernova remnant (G 78.2+2.1). ¹⁷ Pulsations with $P \sim 265$ ms were detected using blind periodicity searches [2] and the pulsar was ¹⁸ associated with the bright EGRET source EG J2020+4017. A likely X-ray counterpart was inferred ¹⁹ using *Chandra* data [3] and X-ray pulsations were later observed with XMM-*Newton* [4]. ²⁰ Among the more than 250 γ -ray pulsars observed by *Fermi*-LAT ¹, PSR J2021+4026 shows a

²¹ unique nature. In fact, its γ -ray flux ($F_{\gamma} \sim 7.9 \times 10^{-10} \,\mathrm{erg \, cm^{-2} \, s^{-1}}$) varies rapidly and simultaneously ²² with its spin-down rate ($\dot{P} \sim 5.4 \times 10^{-14} \,\mathrm{s \, s^{-1}}$) at intervals of a few years, periodically switching ²³ between high-flux/low-spin-down states and low-flux/high-spin-down states. The first state change ²⁴ (*jump*) was observed in October 2011 [5], when the flux dropped by ~ 18% and the spin-down rate ²⁵ increased by ~ 5.6% in a timescale < 7 days. After a slow recovery, occurred over ~ 100 days ²⁶ around December 2014 [6] [7], a similar jump occurred in February 2018 [8]. Previous works also ²⁷ studied the variations of the γ -ray spectrum and of the pulse profile.

We have another example of variability among γ-ray pulsars, which is the binary millisecond pulsar PSR J1023+0038 [9]. In 2013, this source showed an increase in the γ-ray flux concurrent with a radio disappearance [10]. However, this event was modeled as a transition between a rotation-powered state and an accretion-powered state due to the presence of a companion. On the other hand, the jumps of PSR J2021+4026 were modeled as starquake-induced glitches associated to a magnetospheric reconfiguration [6]. Therefore, PSR J2021+4026 is currently the only known gamma-ray variable pulsar.

Here we present the preliminary results of an advanced maximum likelihood analysis based on the most updated *Fermi*-LAT data. Our purpose is to accurately measure variations in the γ -ray flux and spectrum across the jumps. We will discuss the observations in the perspective of a change in the geometry of the pulsar magnetosphere.

39 2. Methods

⁴⁰ Data were prepared and analyzed using *Fermitools*², the standard analysis suite released by ⁴¹ the *Fermi*-LAT collaboration. Our dataset covers 12 years from August 5, 2008 to May 26, 2020, ⁴² and it includes all LAT photons of P8R3_SOURCE_V2 event class. We selected events in a region ⁴³ with a radius of 10° around PSR J2021+4026, with zenith angles $z < 90^{\circ}$ and in the energy range ⁴⁴ from 100 MeV to 300 GeV. We binned data with 35 logarithmically spaced energy bins (10 bins per ⁴⁵ decade) and squared angular pixels of size 0°.1. Figure 1 shows the produced counts map. ⁴⁶ LAT photons are partitioned into four event types (PSF0, PSF1, PSF2, PSF3) based on the

LA1 photons are partitioned into four event types (PSF0, PSF1, PSF2, PSF3) based on the
 quality of the angular reconstruction. Each event type has a different LAT response; therefore, we
 prepared four sets of binned data and performed a maximum likelihood analysis with summed PSF
 components. In order to study the spectral variations between different states, we further divided
 data in four distinct time intervals, labelled with capital Roman letters. These intervals are defined

¹https://confluence.slac.stanford.edu/display/GLAMCOG/Public+List+of+LAT-Detected+ Gamma-Ray+Pulsars

²https://fermi.gsfc.nasa.gov/ssc/data/analysis/documentation/



Figure 1: Binned counts map including all photons collected between August 5, 2008 and May 26, 2020. The size of the RoI is $14^{\circ} \times 14^{\circ}$. The size of the pixels is 0° .1. The three bright sources are PSR J2021+4026 (center), PSR J2021+3651 (3° .6 south) and PSR J2032+4027 (2° .3 north-east). The Galactic plane is visible in the background.

as follows. A: August 5, 2008 (MJD 54683) - October 16, 2011 (MJD 55850). B: October 16, 2011

⁵² (MJD 55850) - December 9, 2014 (57000). C: December 9, 2014 (MJD 57000) - February 2, 2018

⁵³ (MJD 58150). D: February 2, 2018 (MJD 58150) - May 26, 2020 (MJD 58995).

⁵⁴ We built a model of the γ -ray sky starting from the 4FGL catalog [11] and including all sources ⁵⁵ within 20° from PSR J2021+4026. We also included templates for the Cygnus Loop and for the ⁵⁶ Galactic and isotropic diffuse emissions. The γ -ray spectrum of PSR J2021+4026 was modeled as ⁵⁷ a power law with an exponential cutoff,

$$dN/dE \propto E^{-\Gamma_1} \exp\left(-bE^{2/3}\right) \quad , \tag{1}$$

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where the normalization and the spectral parameters were kept free. We freed the normalization of other bright pulsars (PSR J2021+3651, PSR J2032+4027) and extended sources (SNR G 78.2+2.1, Cygnus Cocoon). No other pulsar in the field show variability; however, there are extragalactic sources with reported flaring behavior. In order to take account of spurious flux changes, we freed the normalization of all variable point sources within 7°. Finally, we freed the Galactic diffuse emission and fixed the isotropic diffuse emission. We ended up with 23 free parameters. In the binned likelihood analysis, the accuracy on the results at energies below 1 GeV was

⁶⁵ improved by applying corrections due to the energy dispersion. Moreover, we took account of
 ⁶⁶ systematic errors on the diffuse background by including likelihood weights, which were calculated
 ⁶⁷ for each event type.



Figure 2: Fitted spectral energy distributions of PSR J2021+4026 in intervals A (red), B (green), C (blue) and D (orange) at the 2011 and 2018 jumps. The bands represent the 3σ credibility intervals from a multivariate Gaussian distribution. The inset panels show the 3σ credibility ellipses around the optimal values of power-law index, Γ_1 , and exponential factor, *b*.

68 **3.** Results

The fits to the time intervals A, B, C, and D show consistent relative variations in the γ -ray flux 69 at both state changes: $\Delta F_{\gamma}/F_{\gamma} = -15.3 \pm 1.4\%$ at the 2011 jump, $\Delta F_{\gamma}/F_{\gamma} = -16.3 \pm 1.4\%$ at the 2018 70 jump. We can also observe differences in the spectrum, mainly due to variations in the exponential 71 factor, b, rather that the power-law index, Γ_1 , which does not change significantly. We also ran a fit 72 using a model with fixed spectral parameters (*global*). The significance of the variable model with 73 respect to the global model is > 3σ in A, B and D, while significance is lower (2σ) in C. However, 74 the pre-jump and post-jump spectra are significantly different (3σ) at both jumps (Figure 2). 75 We analyzed 30-day and 7-day time intervals in order to study short-timescale variations. Due 76 to the reduced exposure, we were only able to fit the gamma-ray flux of PSR J2021+4026; therefore, 77 we fixed the spectral parameters in each state to the values from the global model. We obtained 78 estimates of the relative variations at the jumps starting from the mean values and the standard 79 deviations in intervals A, B, C and D. In particular, we measured $\Delta F_{\gamma}/F_{\gamma} = -16\pm6\%$ in 2011,

⁸⁰ deviations in intervals A, B, C and ⁸¹ $\Delta F_{\gamma}/F_{\gamma} = -13\pm5\%$ in 2018.

We also studied the timing variability by performing an *H*-test [12] on time intervals of 60 days. This way we could relate flux jumps to the corresponding changes in frequency, *f*, and spin-down rate, \dot{f} (Figure 3). We obtained $\Delta \dot{f}/\dot{f} = 5\pm 3\%$ at the 2011 jump, $\Delta \dot{f}/\dot{f} = 5\pm 2\%$ at the 2018 jump.

4. Discussion and conclusions

According to a previously proposed model [6], the changes in spin-down rate and luminosity of the pulsar could be produced by a shift in the magnetic inclination angle, α . Such shift may be the consequence of a glitch induced by a starquake.



Figure 3: Fitted energy flux and optimal timing parameters of PSR J2021+4026 in the time range from August 5, 2008 to May 26, 2020. Rather than the frequency, f - k·MJD is reported, where $k = 6.844 \times 10^{-8}$ Hz day⁻¹ is a global spin-down rate obtained from a χ^2 fit to the *f* data. Horizontal dashed lines and shaded bands represent the mean values and the 1 σ confidence bands in the time intervals A (red), B (green), C (blue) and D (orange). Vertical dashed lines indicate the boundaries of the time intervals.

A numerical solution of the equations for magnetohydrodynamics in a force-free magnetosphere
 [13] gives a formula for the γ-ray spin-down luminosity:

$$L_{\rm sd} \propto f^4 \ (1 + \sin^2 \alpha) \tag{2}$$

⁹¹ By setting $L_{sd} = \dot{E}_{rot}$ we get

$$\Delta \dot{f} / \dot{f} = \sin 2\alpha \ \Delta \alpha \ (1 + \sin^2 \alpha)^{-1} \tag{3}$$

If we assume a pre-jump inclination angle $\alpha = 63^{\circ}$, obtained from a model of the magnetosphere [6], and the spin-down variation we measured, $\Delta \dot{f} / \dot{f} = 5 \pm 2\%$, we get $\Delta \alpha = 6 \pm 2^{\circ}$ and $\Delta L_{sd} / L_{sd} = 5.0 \pm 1.7\%$. The latter value is not consistent with the measured relative flux variations, indicating that this simple model is not sufficient to describe the behavior of PSR J2021+4026. Further information may be obtained from a detailed study of the pulse profile, which could strengthen or disprove this hypothesis. ⁹⁸ More sophisticated models [14] are able to produce detailed magnetospheric structures; how-⁹⁹ ever, they are valid only in the stationary case. For this reason, we must search hints about the ¹⁰⁰ dynamics of the variations by fitting these models to the observed γ -ray emission in the different ¹⁰¹ states. In this context, the variability analysis of γ -ray pulsars has a major importance.

¹⁰² Due to its uniqueness, PSR J2021+4026 plays a key role in pulsar physics. Therefore, we are ¹⁰³ monitoring and analyzing its γ -ray flux continually. This investigation may represent the starting ¹⁰⁴ point of a fine search for variable γ -ray pulsars. Hints of flux variability in other γ -ray pulsars have ¹⁰⁵ already appeared in the 4FGL catalog, as the case of PSR J2043+2740 [11]. For this reason, we ¹⁰⁶ believe that this approach will produce significant results and will lead to a deeper knowledge of ¹⁰⁷ the physics of pulsar magnetospheres.

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