Hα polarimetry as a powerful diagnostics of cosmic-ray modified shock

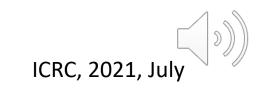
Cf.) Shimoda & Laming 2019a, MNRAS, 485 Shimoda & Laming 2019b, MNRAS, 489

<u>Jiro Shimoda¹</u>

J. Martin Laming²



1. Nagoya Univ.; 2. Naval Research Lab.



Summary

- Cosmic-Ray Modified Shocks (CRMSs) are one of an essential prediction of the diffusive shock acceleration.
- \square We must examine a velocity modification of plasma with ~ 10 % level around the SNR shock.
- The polarization direction of H α responds *sensitively* whether the shock is modified.



Supernova Remnant (SNR)



γ-ray: electron or proton?

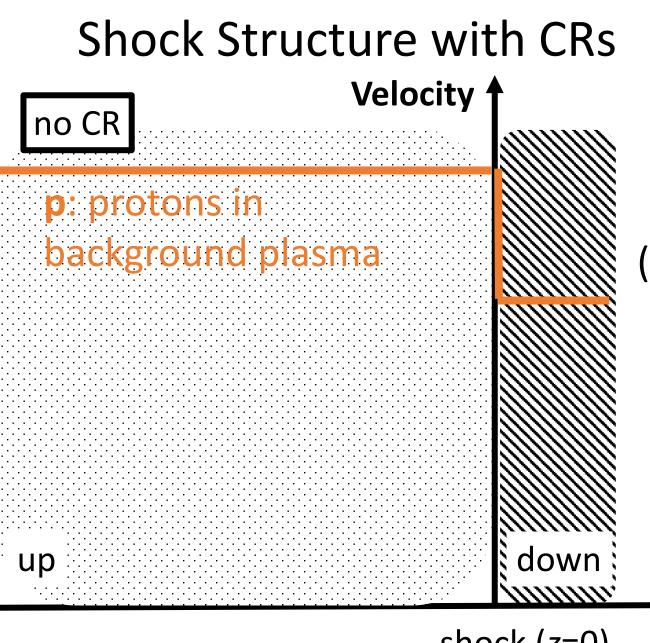
X-ray: ~TeV CR electrons Supernova ejecta

H α : useful tracer of shock condition & physics.

SNR 0509-67.5 (Chandra & HST)

Blue: 1.5 – 7.0 keV Green: 0.2 – 1.5 keV Red: Hα

SNR shock is considered as the best candidate of CR origin.

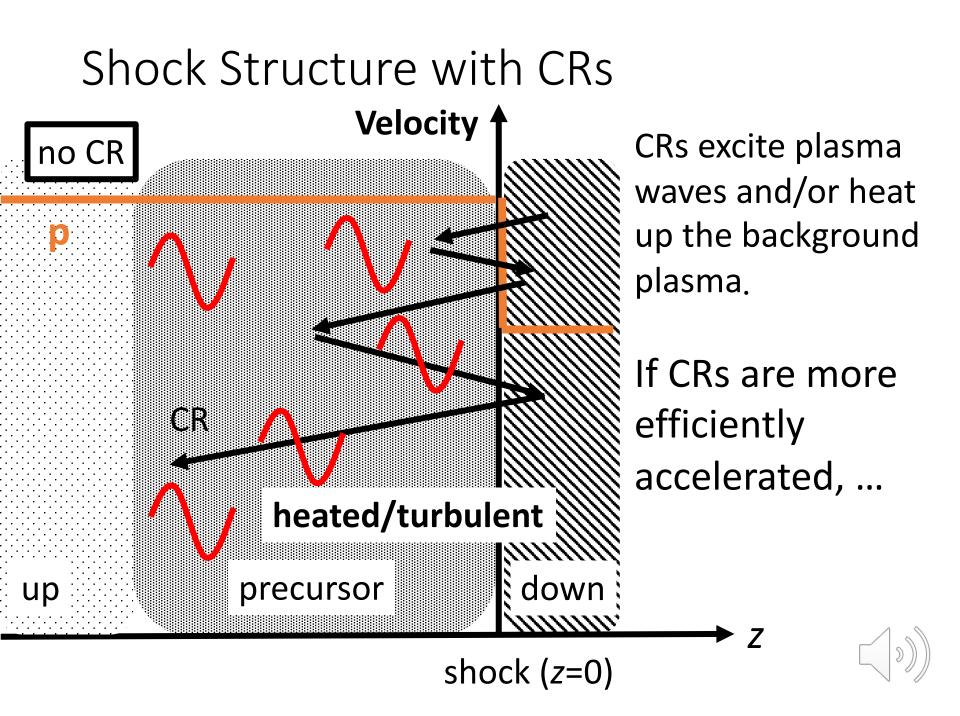


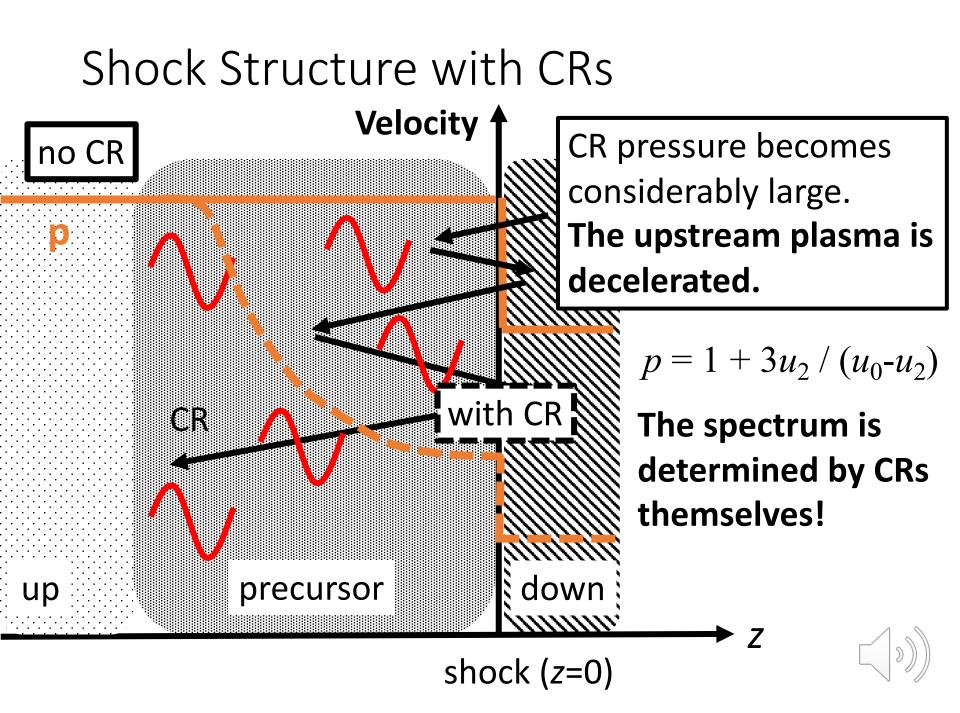
Shock structure for no CRs (Hydrodynamics)

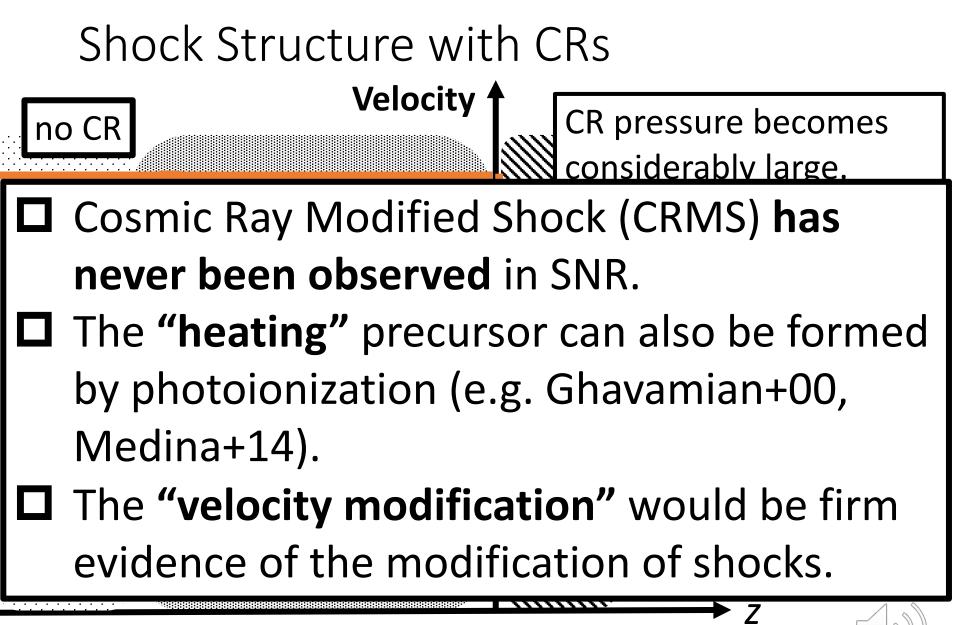
When CRs are accelerated ?

7

shock (z=0)

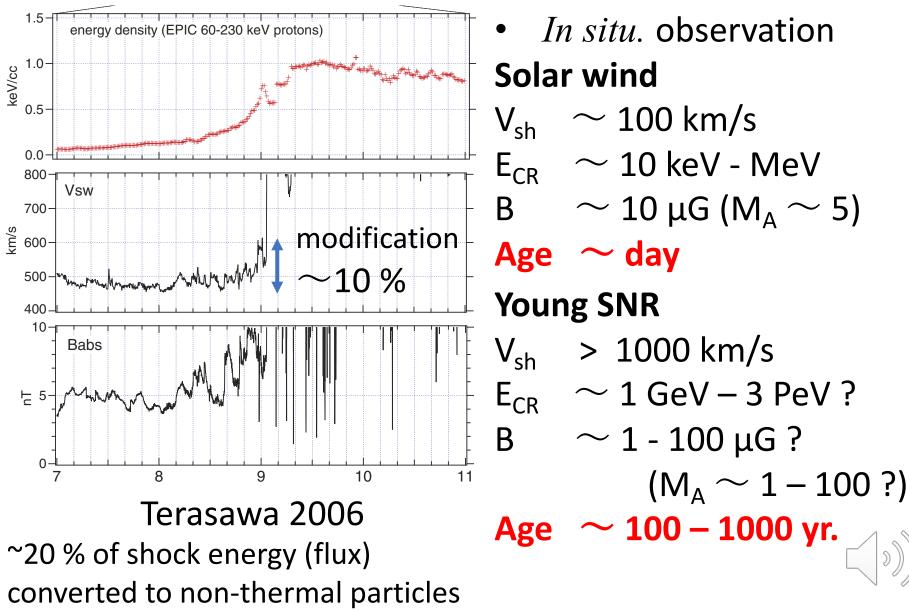




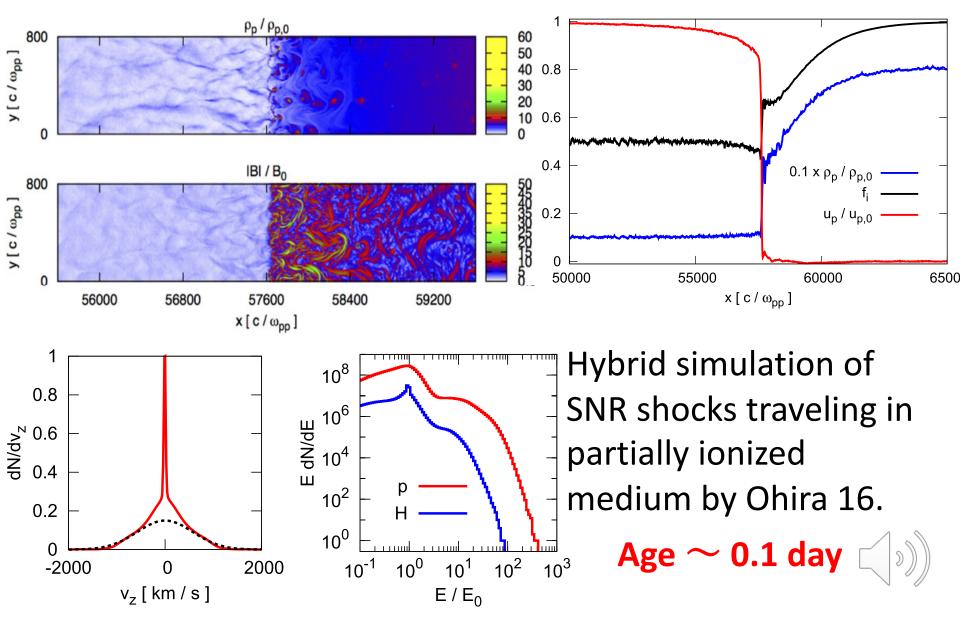


shock (z=0)

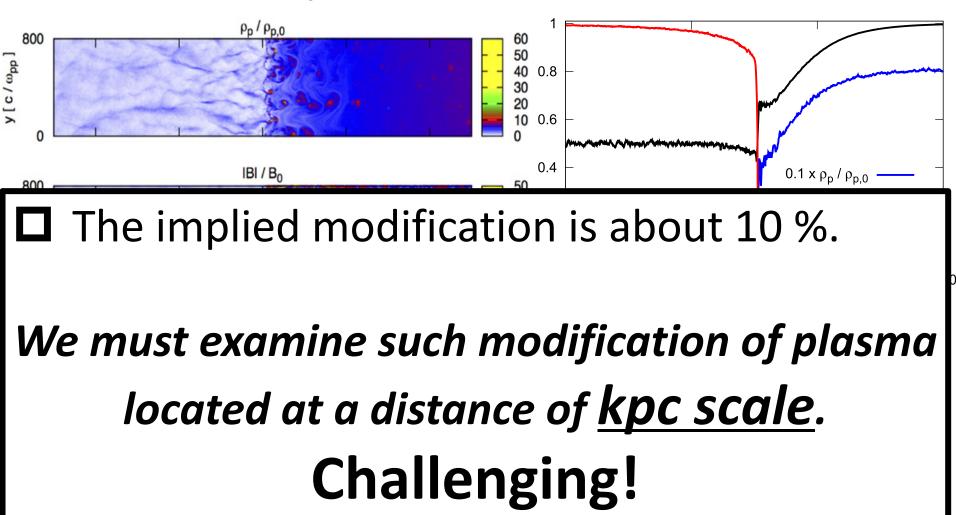
Cosmic-Ray Modified Shock (CRMS)

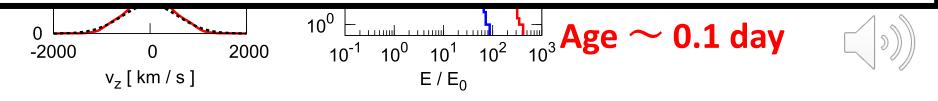


Cosmic-Ray Modified Shock (CRMS)

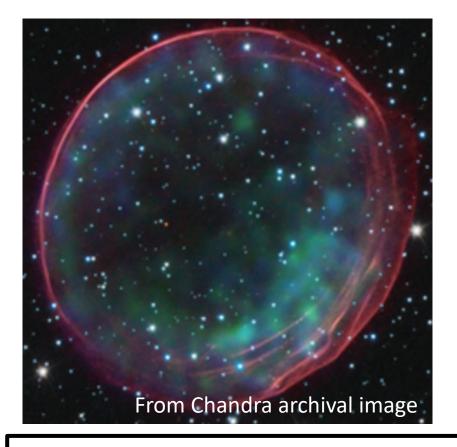


Cosmic-Ray Modified Shock (CRMS)





Supernova Remnant (SNR)



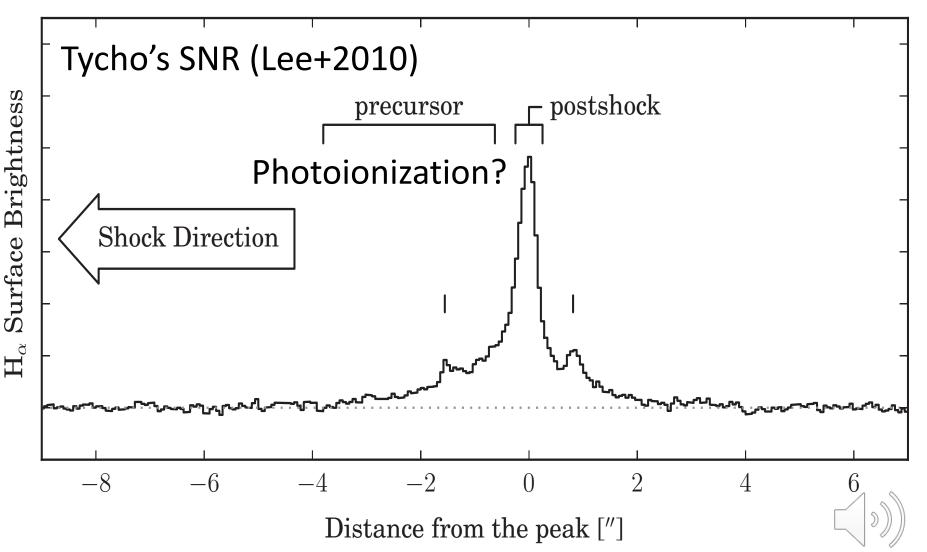
γ-ray: electron or proton?

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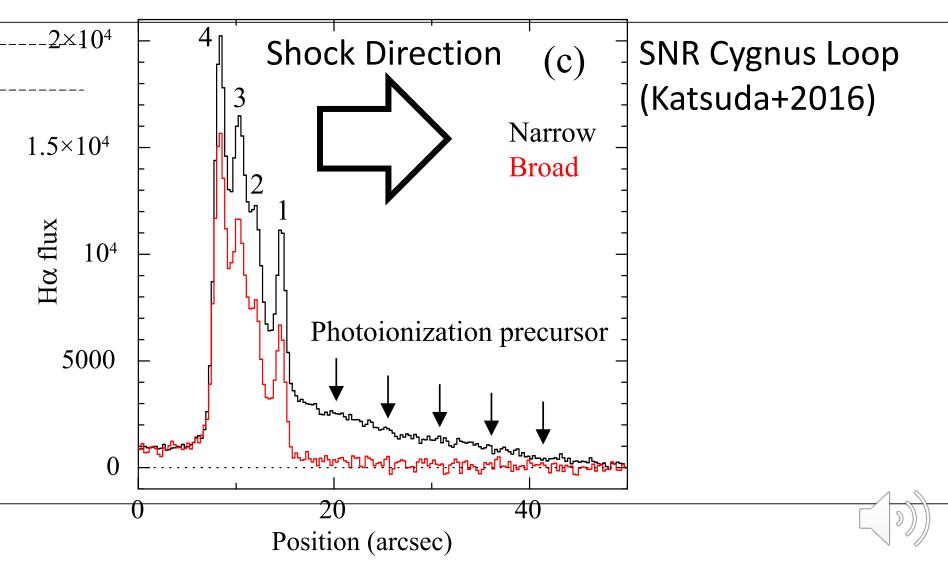
H α : useful tracer of shock condition & physics.

Hα emissions reflect a plasma condition around the shock (e.g. Raymond 91 for review). Red: Hα **best candidate of CR origin**.

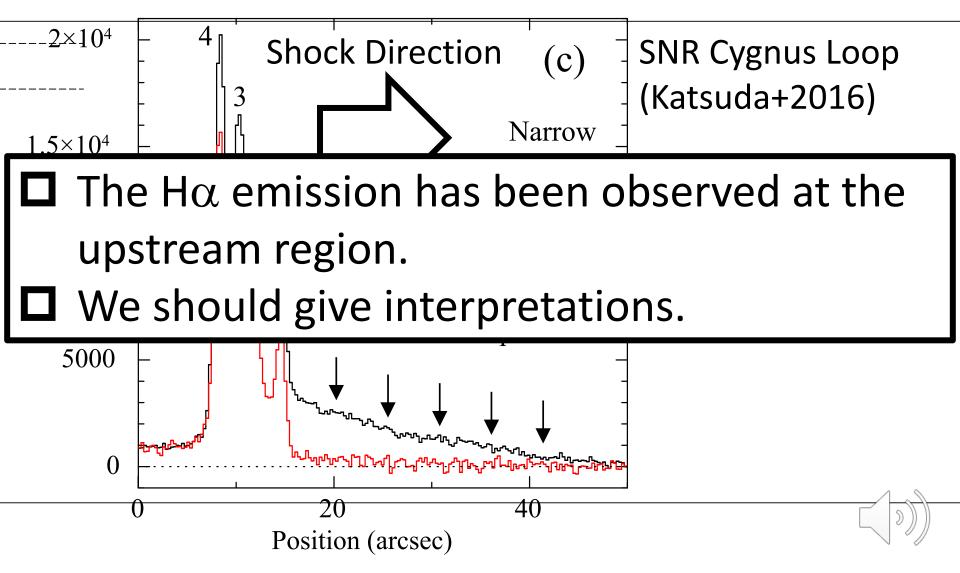
$H\alpha$ emission from upstream

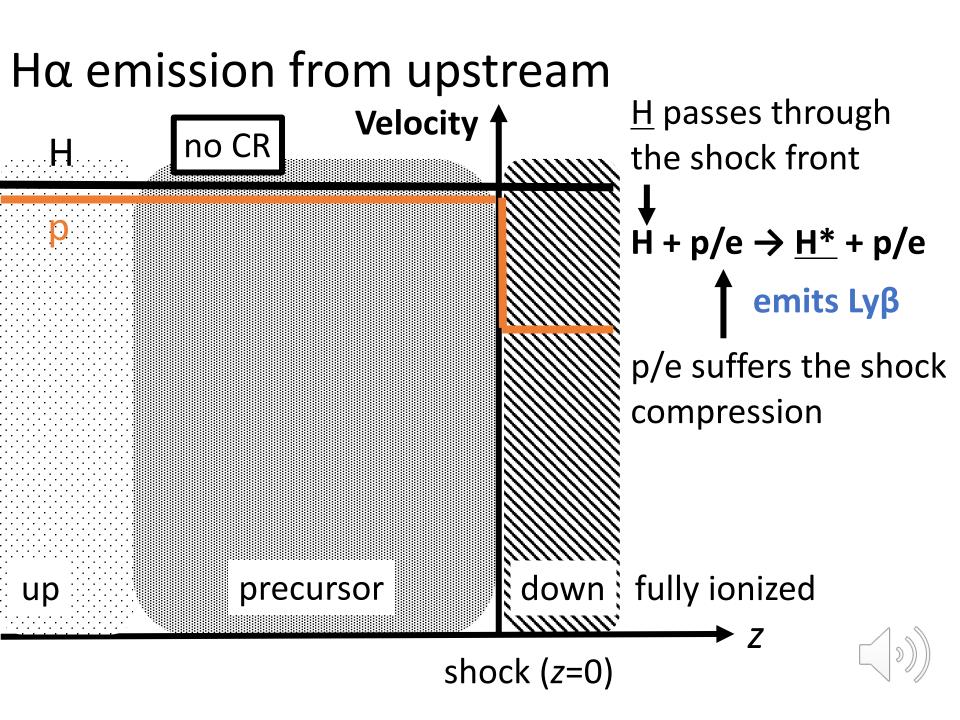


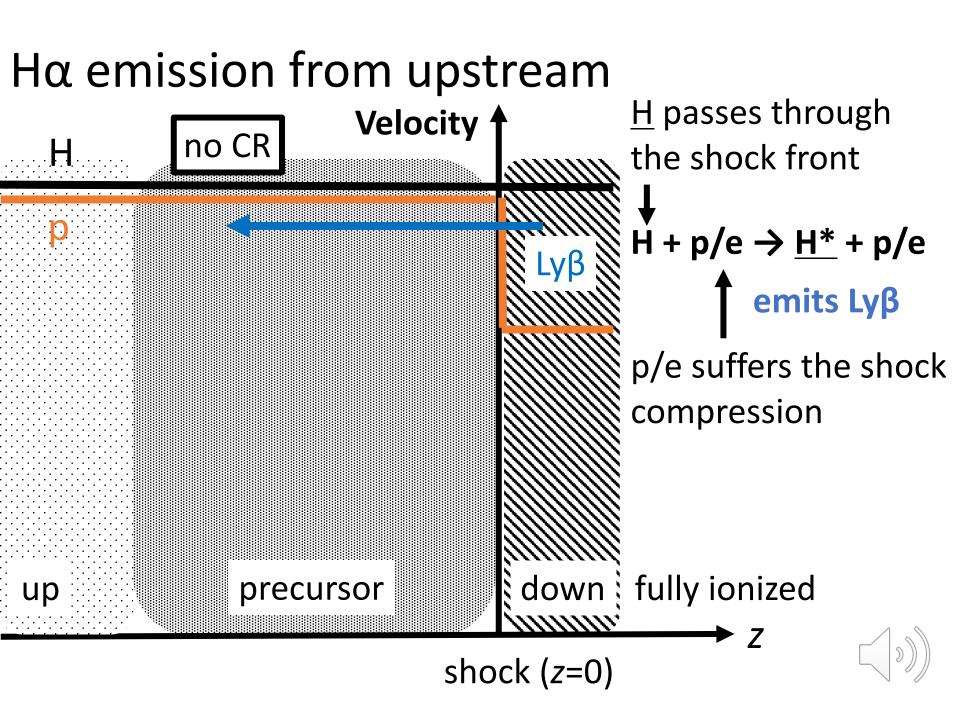
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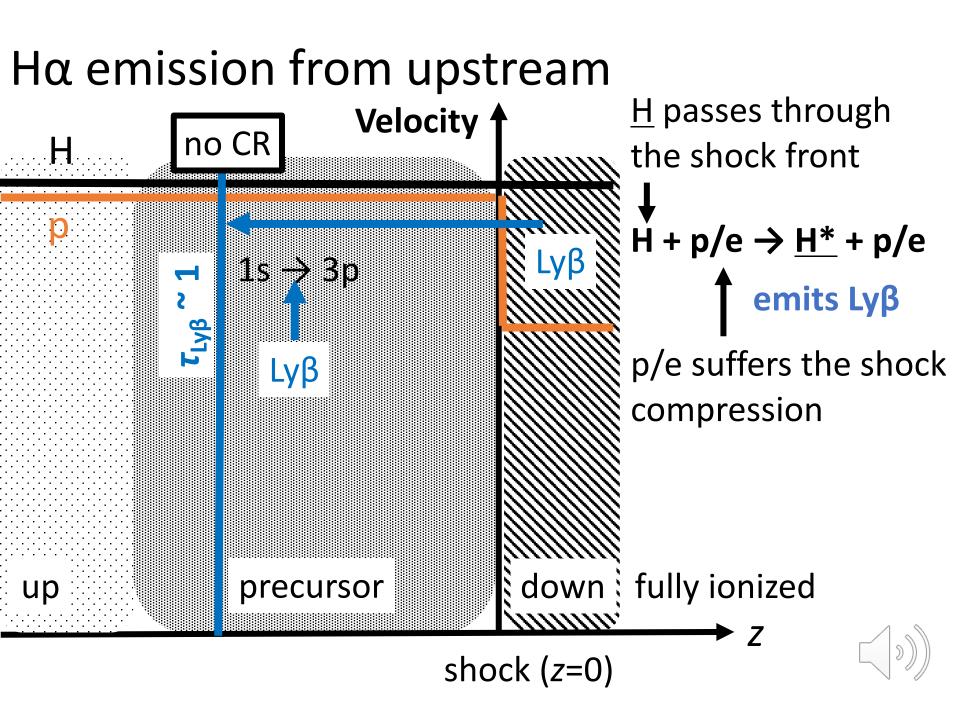


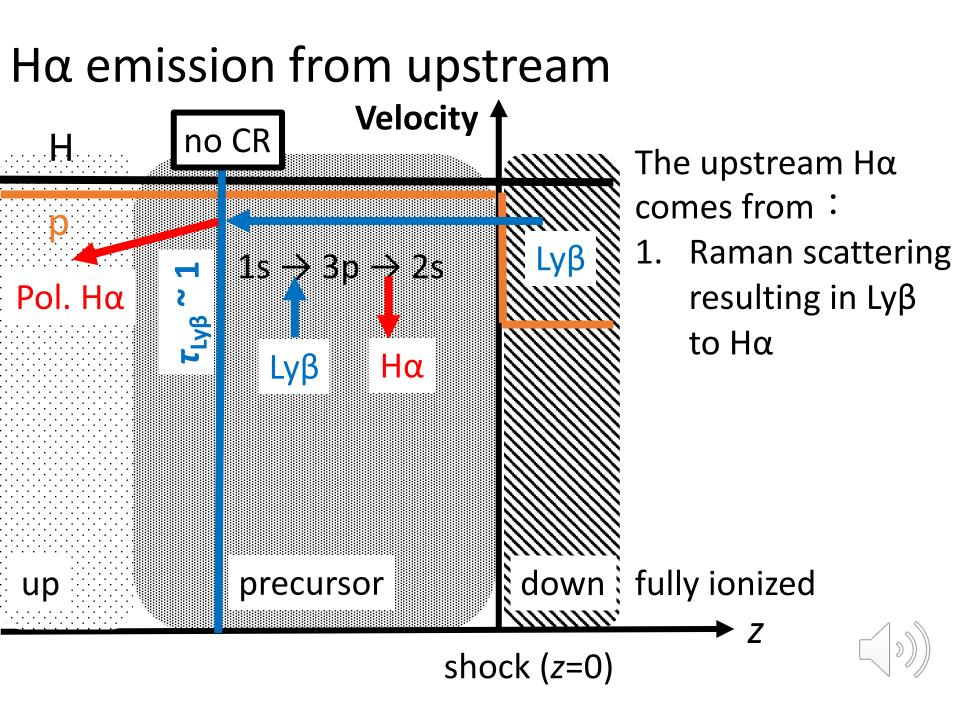
$H\alpha$ emission from upstream

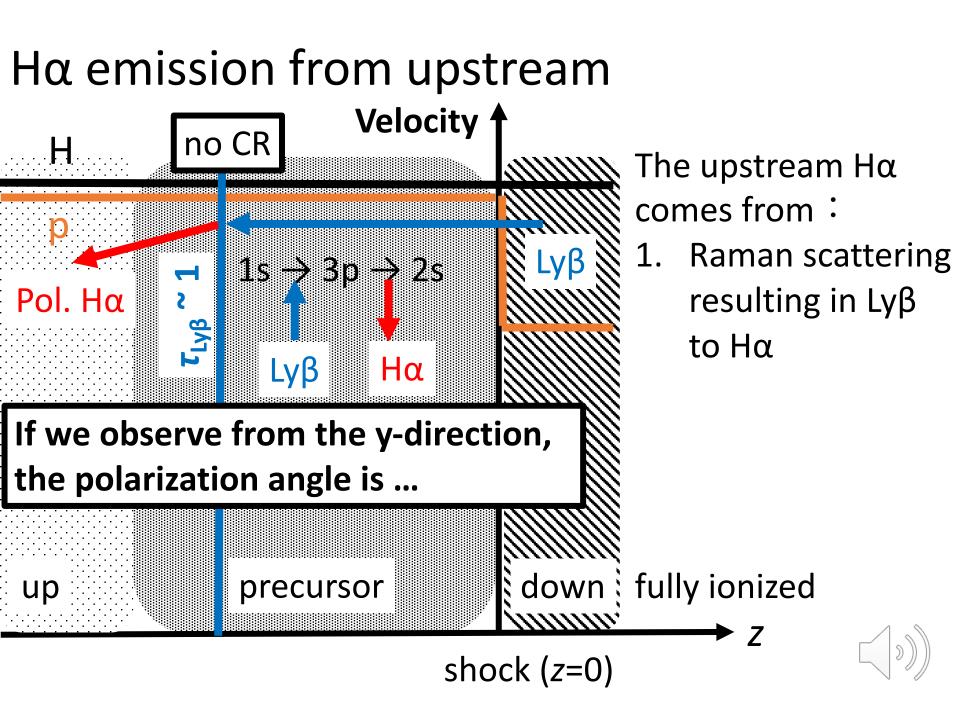




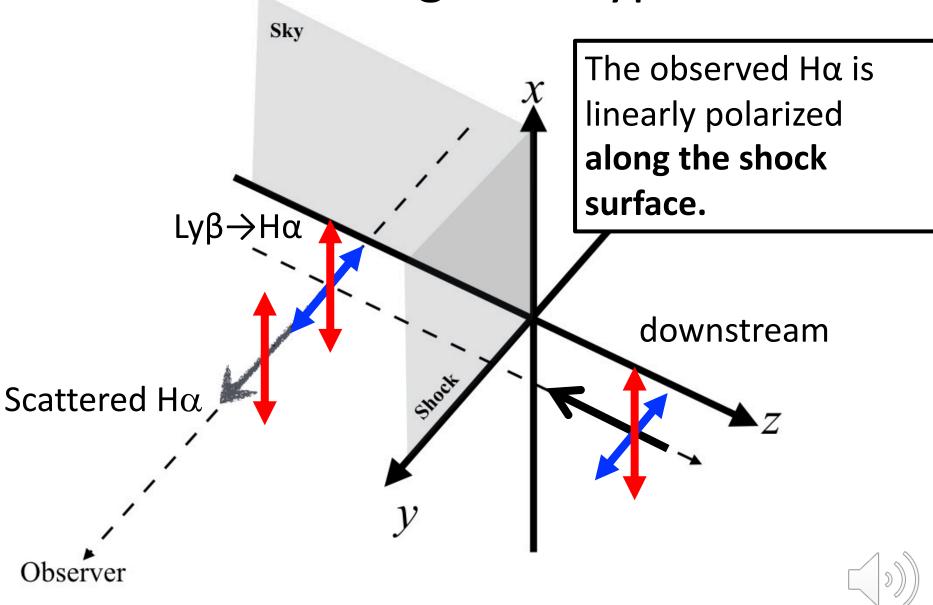


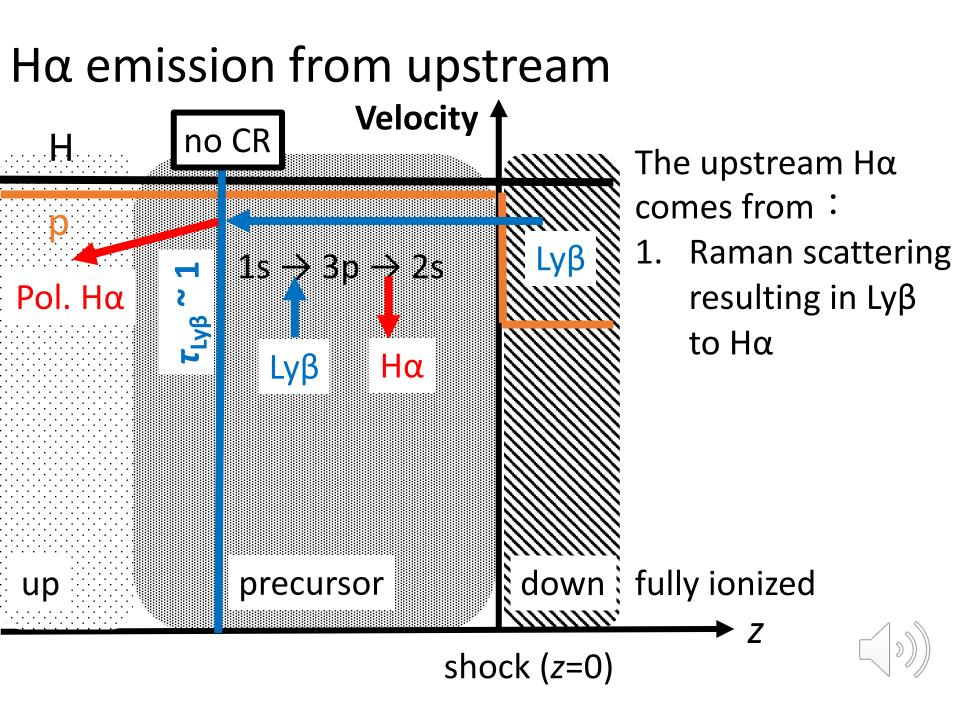


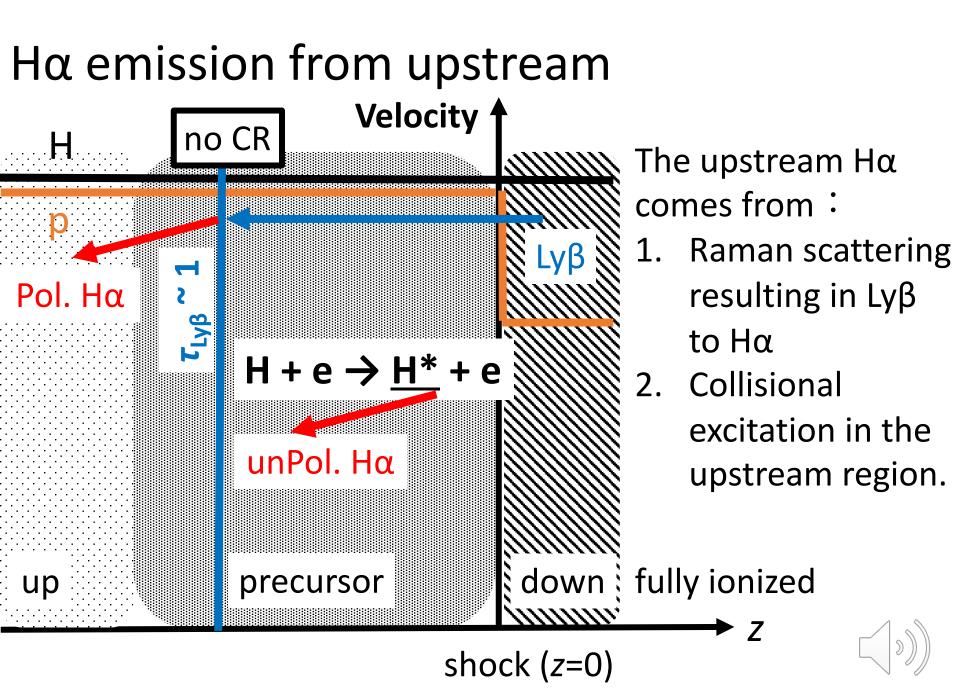


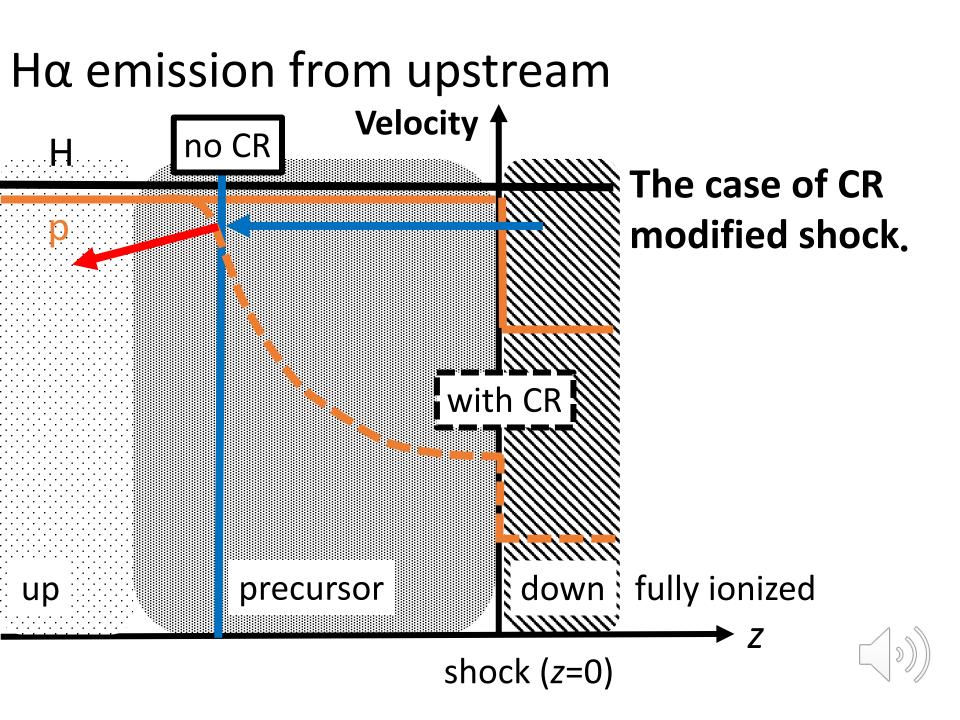


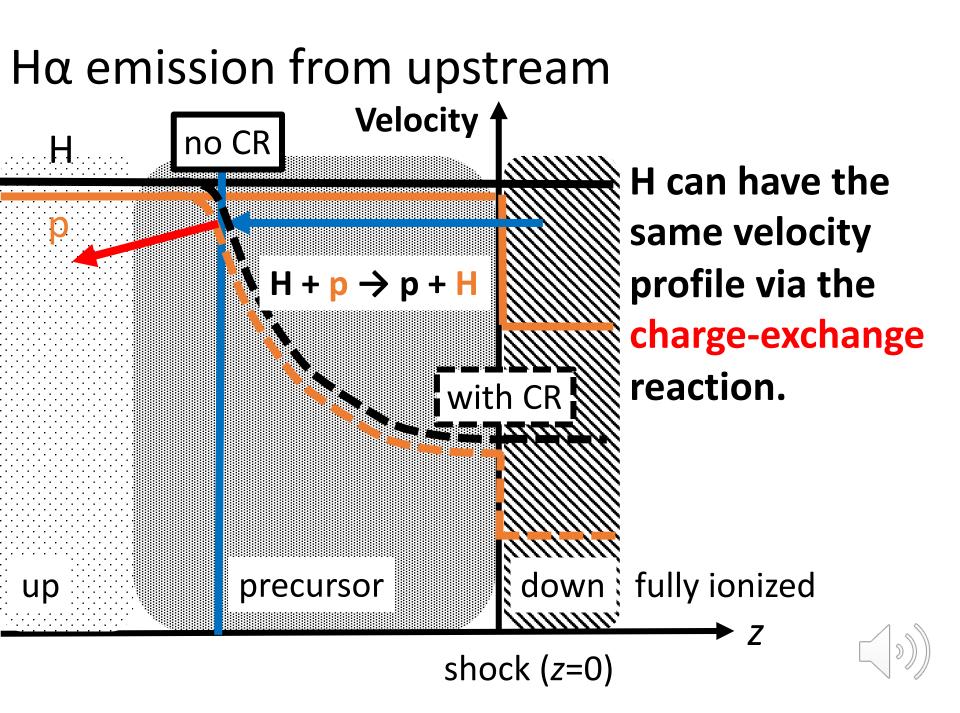
Polarization angle for Ly $\beta \rightarrow H\alpha$

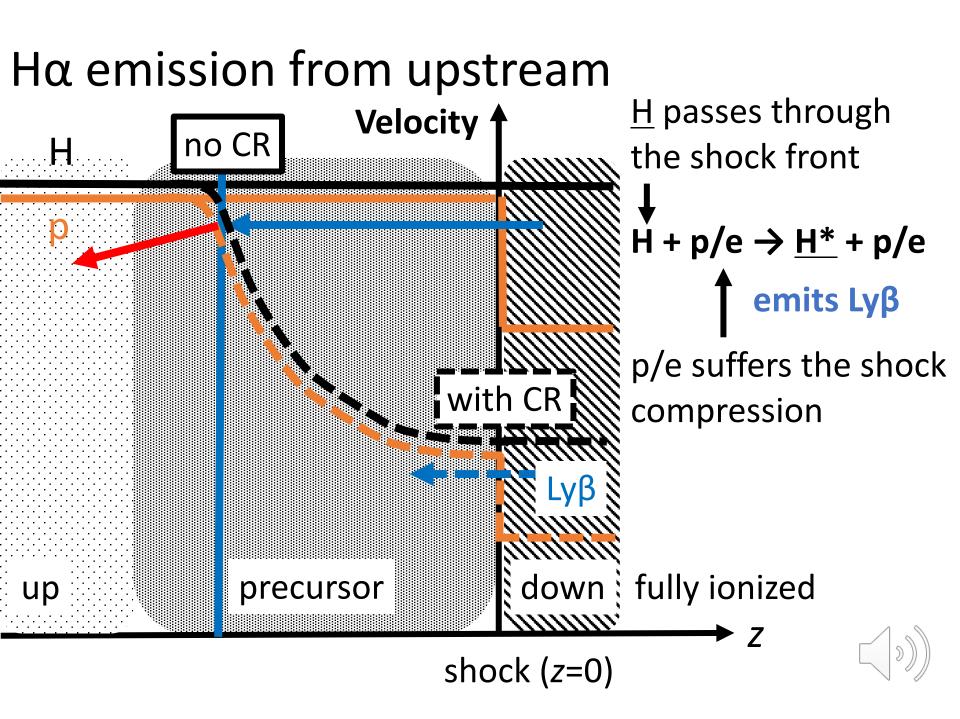


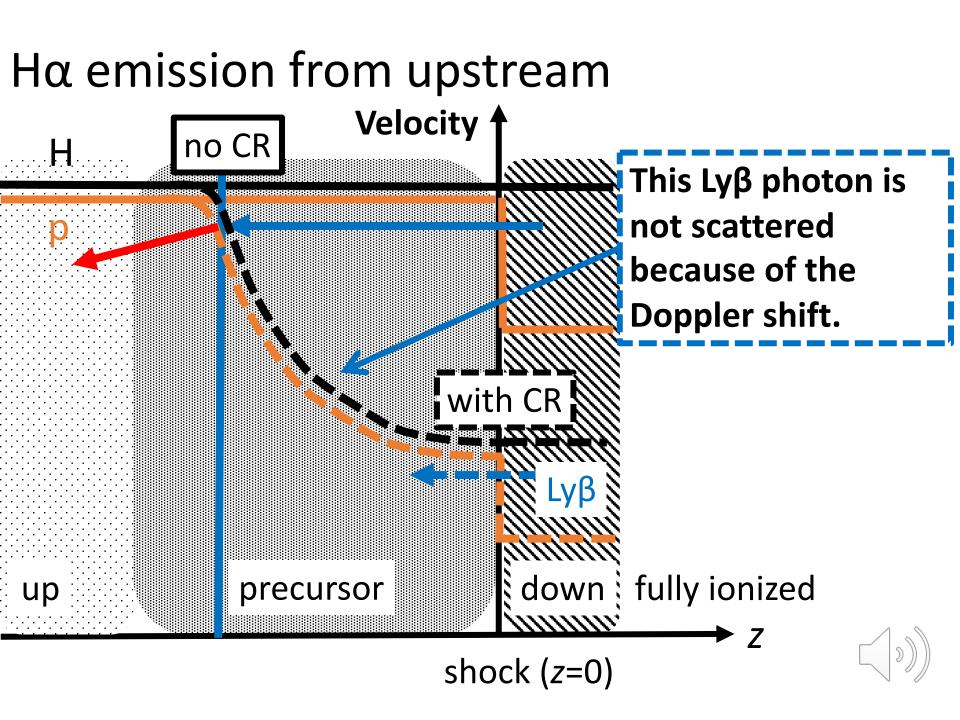


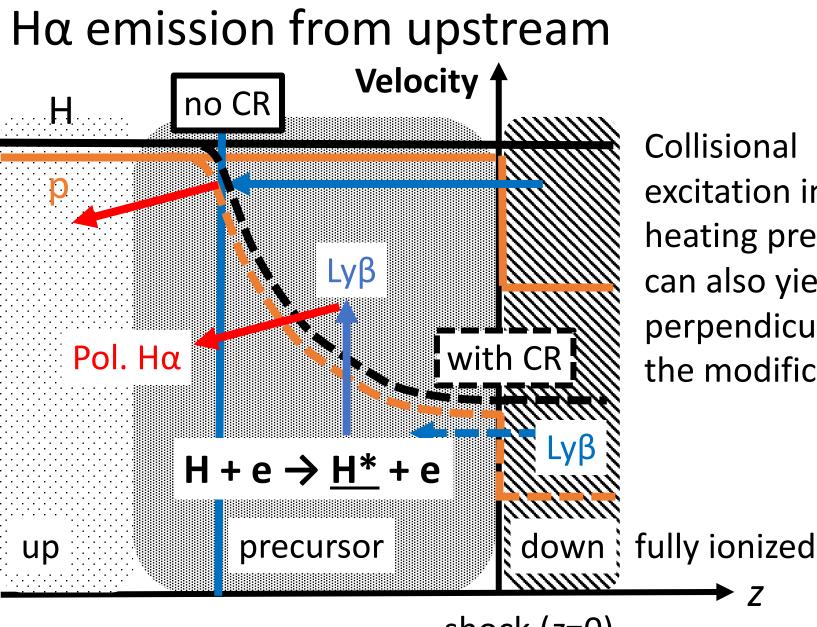








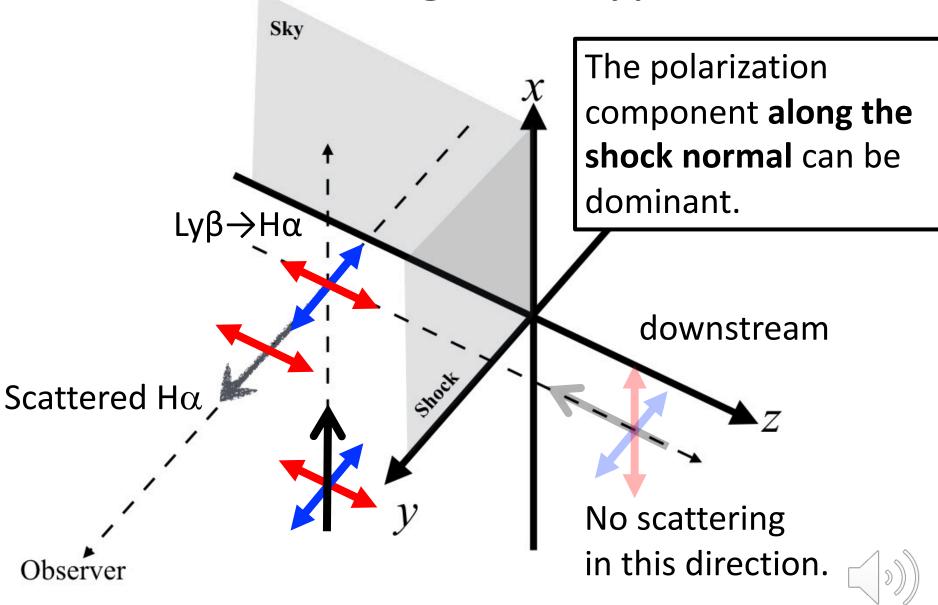


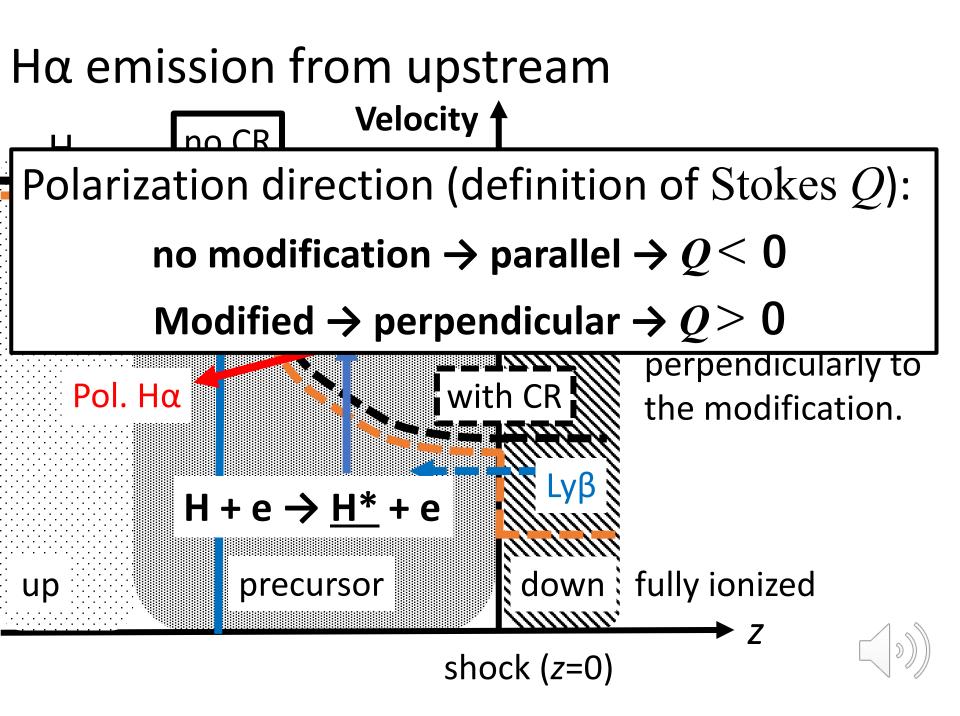


excitation in the heating precursor can also yield Lyβ perpendicularly to the modification.

shock (z=0)

Polarization angle for Ly $\beta \rightarrow H\alpha$

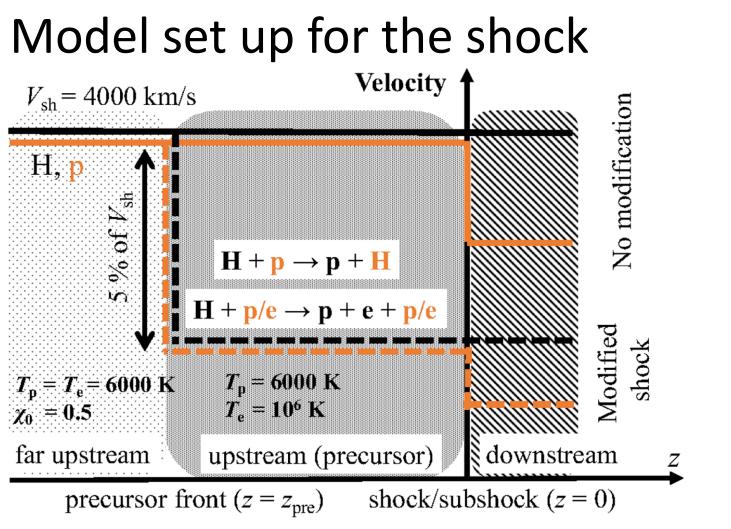




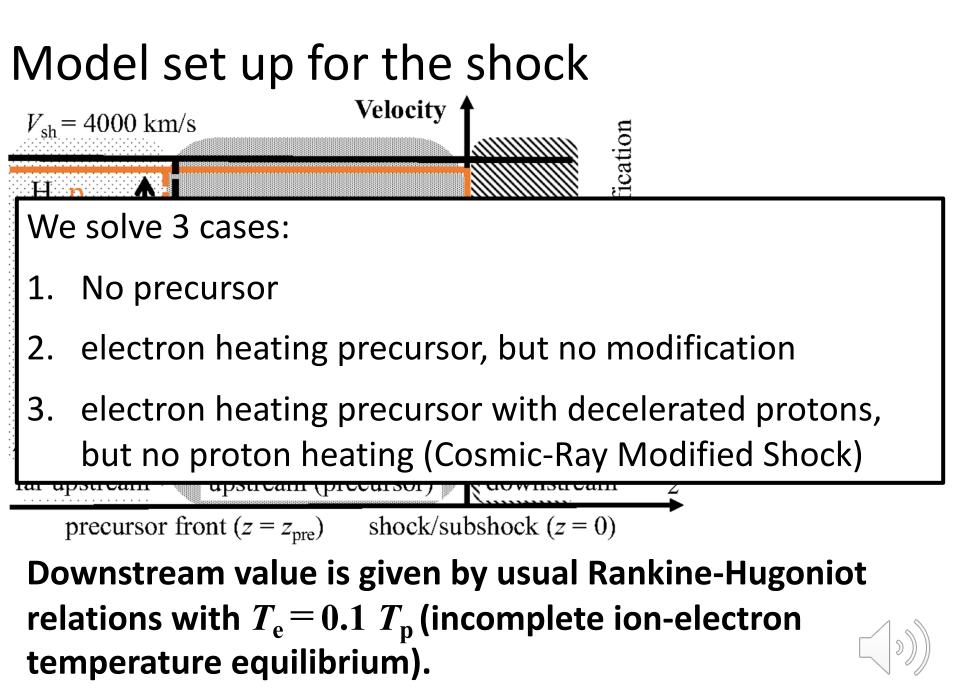
- □ Radiation line transfer & atomic population with polarized light → quite complex!
 □ We make simplifications:
- 1. Omitting the polarization in atomic population calculations (SJ & Laming 19a). Stokes *I* is OK.
- 2. Completely unpolarized Lyβ is supposed.
- 3. For the Stokes *Q*, the 3p_{3/2} state only results from the radiative excitation in the upstream region.

See, SJ & Laming 19b for details

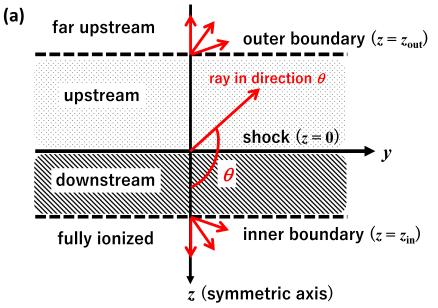


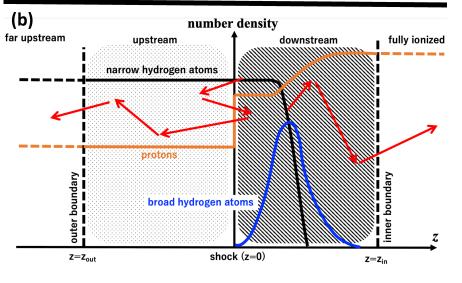


Downstream value is given by usual Rankine-Hugoniot relations with $T_e = 0.1 T_p$ (incomplete ion-electron temperature equilibrium).



Line Transfer Model





Parameters:

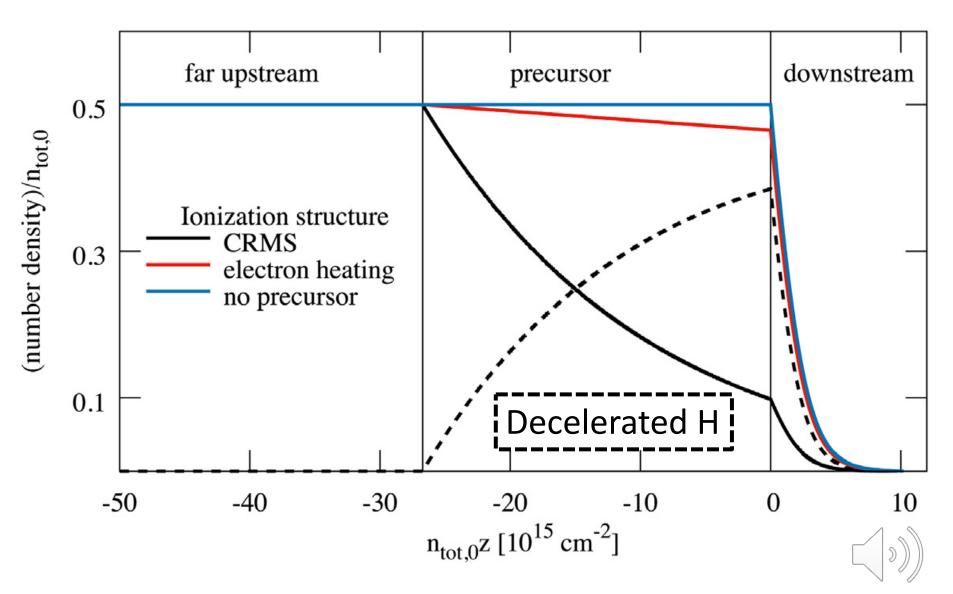
- 1) Shock velocity $V_{\rm sh}$
- 2 Upstream number density $n_{tot,0}$
- (3) proton fraction χ_0
 - ④ Upstream electron temp
- 5 Downstream electron temp $T_{\rm e} = \beta T_{\rm down}$

Pure hydrogen plasma.

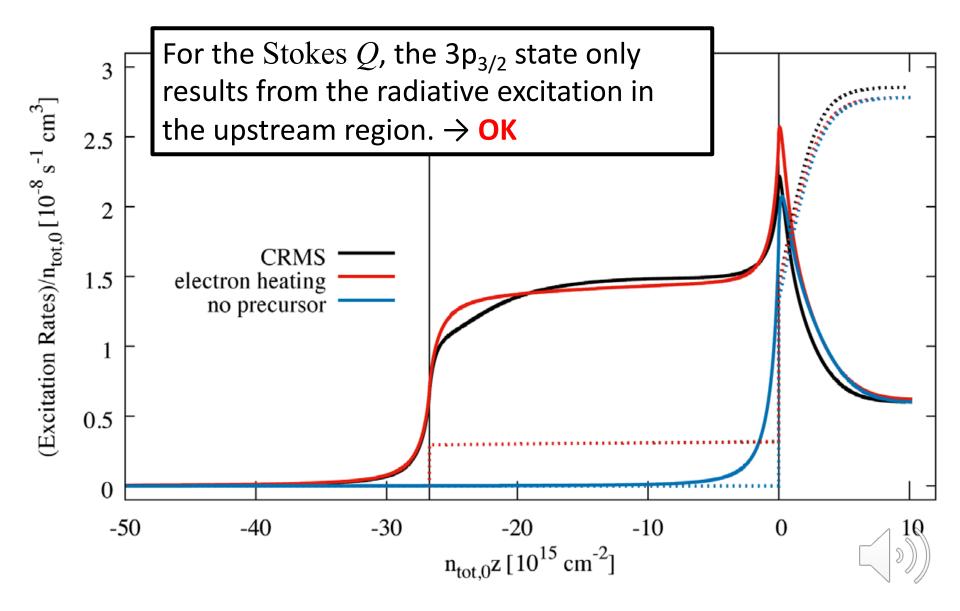
We solve the excited states up to 4f. (SJ & Laming 19a)



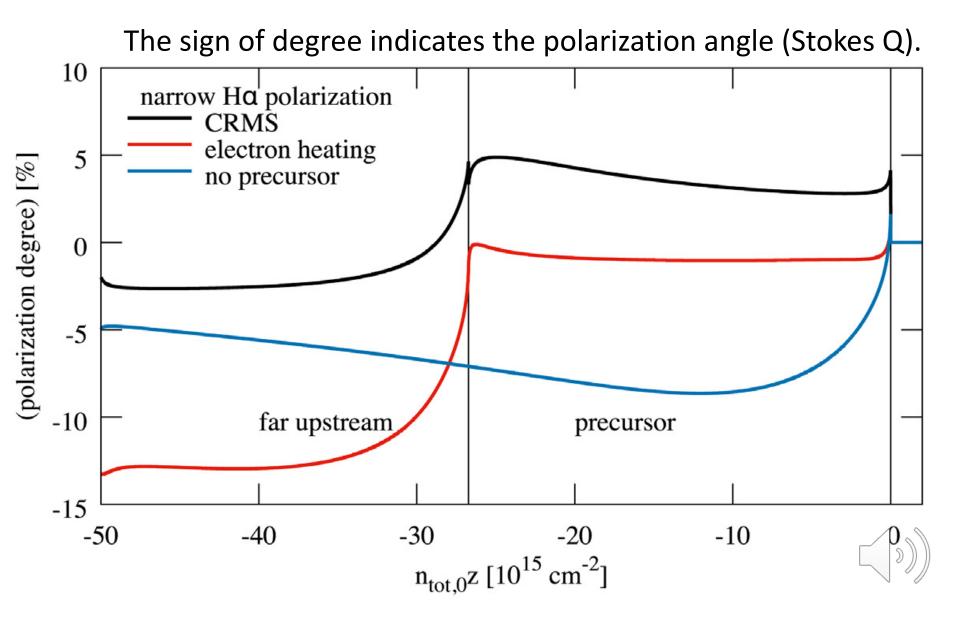
Results: Ionization Structure of H



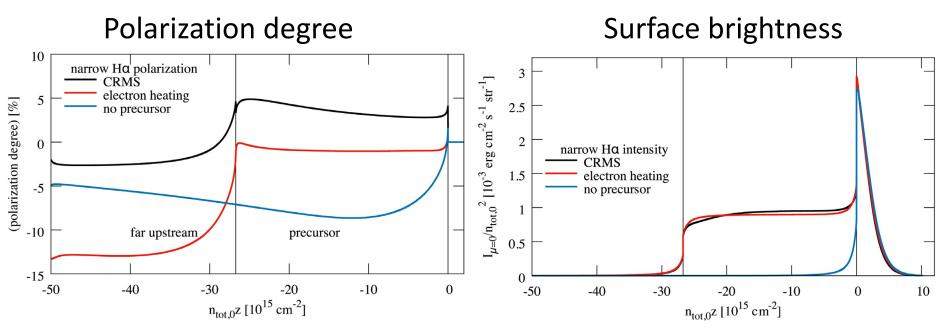
Results: Radiative vs. Collisional



Results: Polarization of $\text{H}\alpha$



Results: Polarization of $\text{H}\alpha$

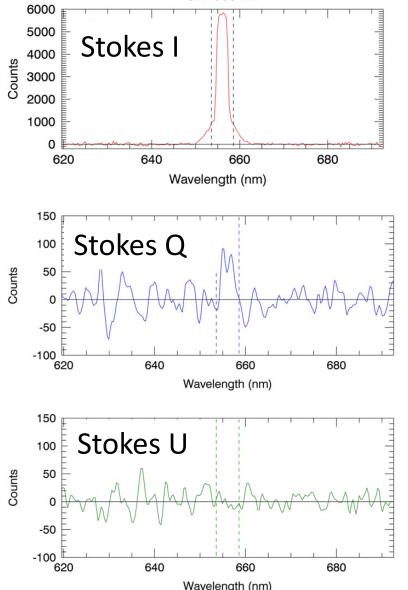


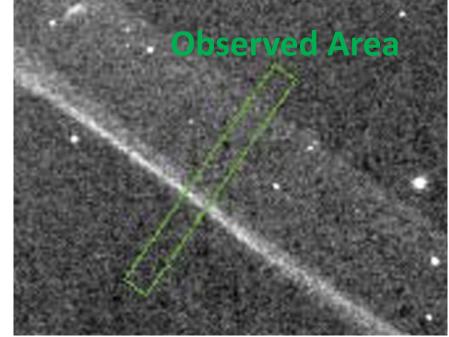
The polarization direction can respond whether the shock is modified.

The degree of a few per cent is measurable (Sparks+ 15).



Discovery of polarized H α emission @ bright filament of SN 1006 (Sparks+ 15)



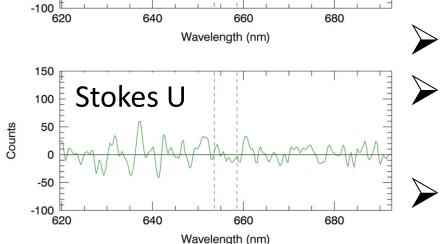


 Linear Polarization
 Polarization angle : perpendicular to the shock
 Degree : 2.0 ± 0.4 %



 \checkmark Our calculation is consistent with the observation.

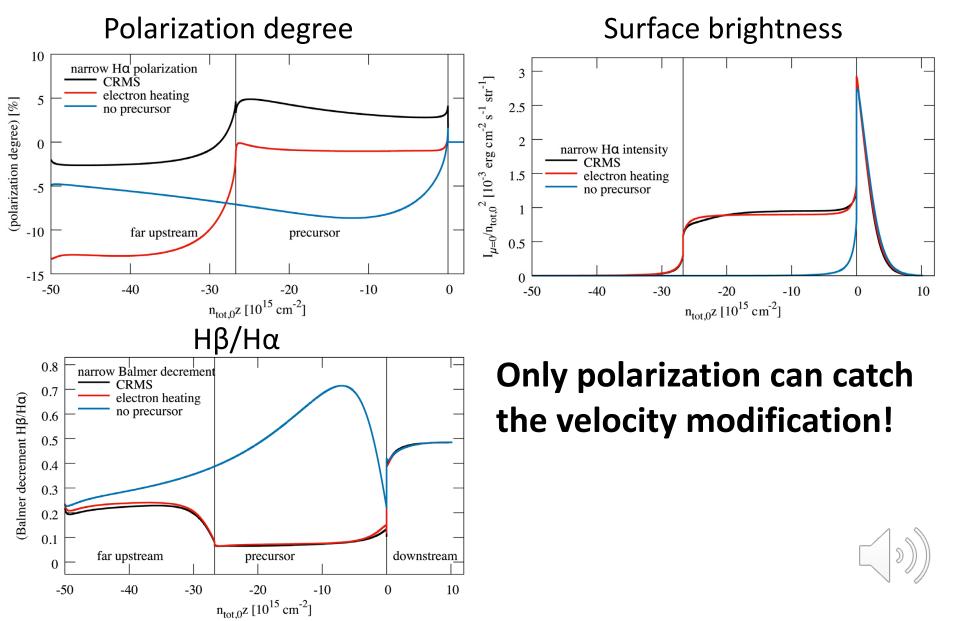
- Polarized Hα has been reported by Sparks+15, but this is not spatially resolved...
- Further observations of Hα polarimetry will bring new insights to particle acceleration!



5000

 Linear Polarization
 Polarization angle : perpendicular to the shock
 Degree : 2.0 ± 0.4 %

Polarization of H α vs others

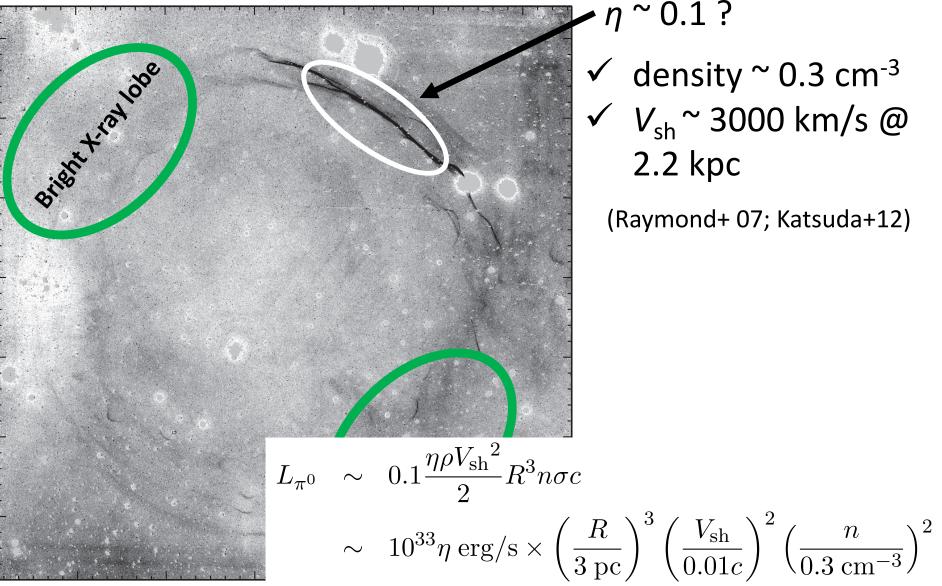


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SN 1006

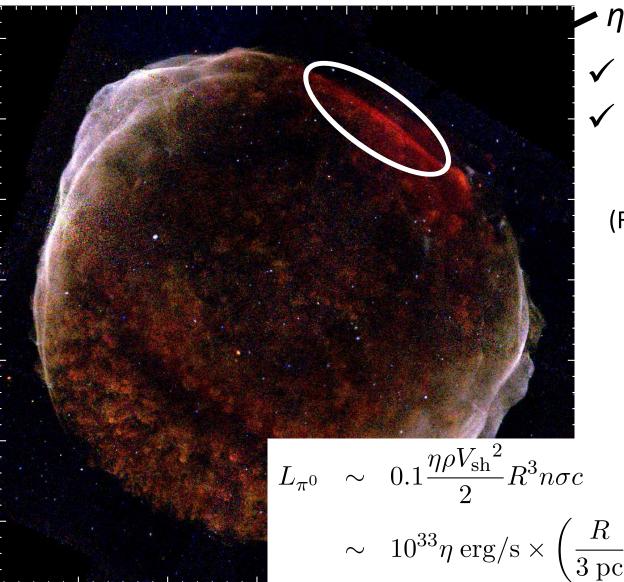


η~0.1?

 \checkmark density ~ 0.3 cm⁻³ ✓ V_{sh}~ 3000 km/s @ 2.2 kpc

(Raymond+ 07; Katsuda+12)

SN 1006

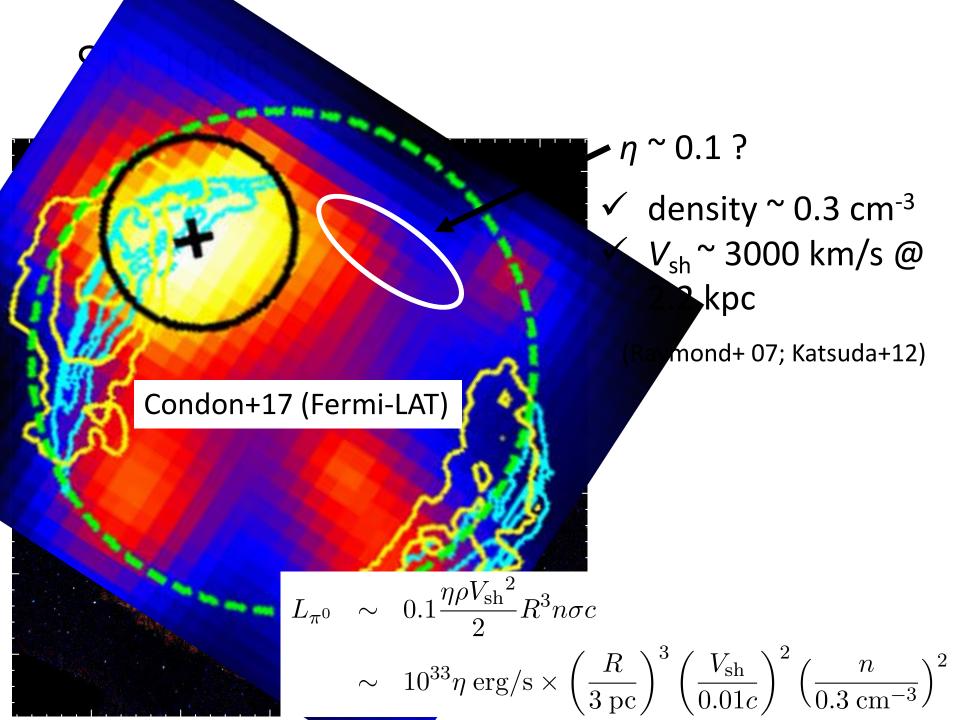


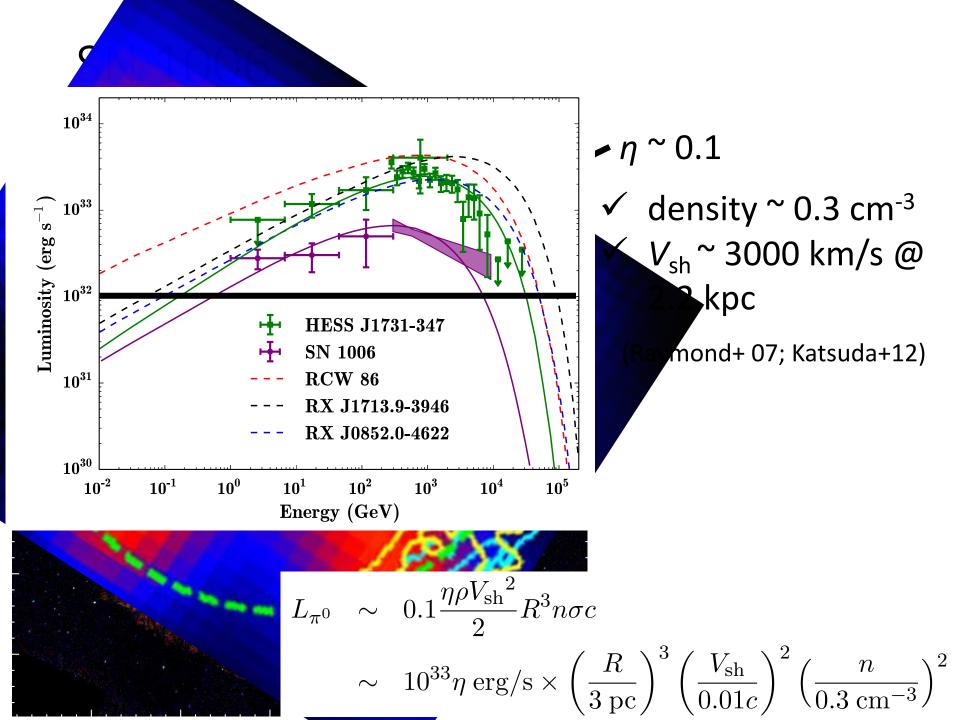
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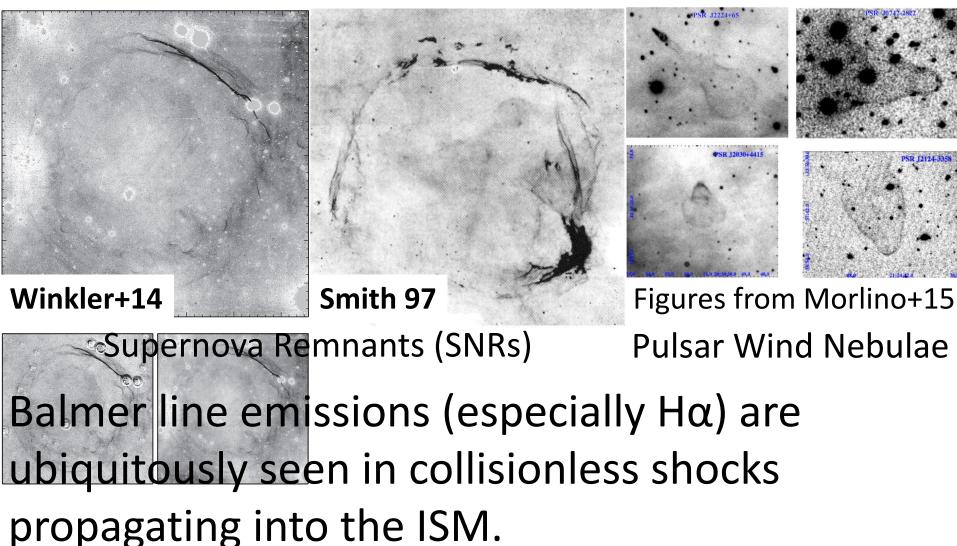
$$\hat{L}_{\pi^0} \sim 0.1 \frac{\eta \rho V_{\rm sh}^2}{2} R^3 n \sigma c$$

 $\sim 10^{33} \eta \, {\rm erg/s} \times \left(\frac{R}{3 \, {\rm pc}}\right)^3 \left(\frac{V_{\rm sh}}{0.01 c}\right)^2 \left(\frac{n}{0.3 \, {\rm cm}^{-3}}\right)^2$

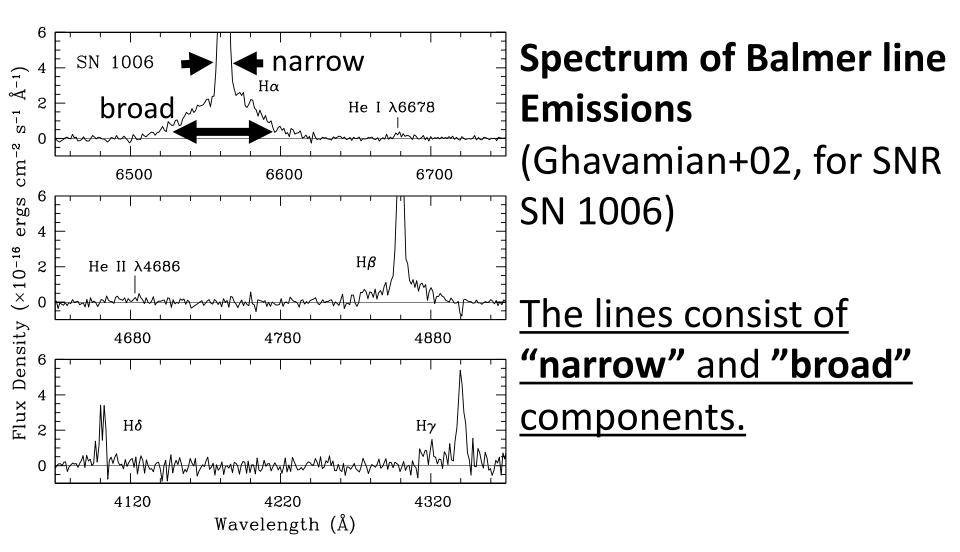




Balmer Line Emissions from Collisionless Shocks



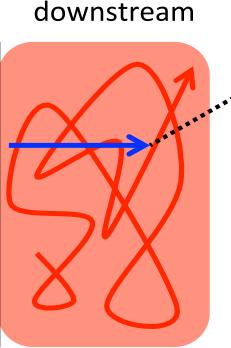
Balmer Line Emissions from Collisionless Shocks



Balmer Line Emissions from Collisionless Shocks

Emission Mechanism (e.g. Chevalier+80)

upstream d



- The collisionless shock is formed by the interaction between <u>charged particles</u> <u>and plasma waves.</u>
- <u>The neutral particles (e.g.</u> <u>hydrogen atoms) are not</u> <u>affected.</u>

SNR shock Charged particles → shock heating Hydrogen atoms → no dissipation

Balmer Line Emissions from Collisionless Shocks Emission Mechanism (e.g. Chevalier+80) downstream upstream Collisional Excitation $H + p (or e) \rightarrow H^* + p (or e)$ Emits "narrow" comp. Charge Transfer $H + p \rightarrow p + H^*$ Emits "broad" comp. The "broad" component reflects the downstream SNR shock Charged particles \rightarrow shock heating temperature of protons. Hydrogen atoms \rightarrow no dissipation