## H $\alpha$ polarimetry as a powerful diagnostics of cosmic-ray modified shock

Cf.) Shimoda \& Laming 2019a, MNRAS, 485 Shimoda \& Laming 2019b, MNRAS, 489

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## Summary

$\square$ Cosmic-Ray Modified Shocks (CRMSs) are one of an essential prediction of the diffusive shock acceleration.
$\square$ We must examine a velocity modification of plasma with $\sim 10$ \% level around the SNR shock.
$\square$ The polarization direction of $\mathrm{H} \alpha$ responds sensitively whether the shock is modified.

## Supernova Remnant (SNR)



## $\gamma$-ray: electron or proton?

X-ray: $\sim \mathrm{TeV}$ CR electrons Supernova ejecta

Ha: useful tracer of shock condition \& physics.

SNR 0509-67.5 (Chandra \& HST)
Blue: $1.5-7.0 \mathrm{keV}$
Green: $0.2-1.5 \mathrm{keV}$ Red: $\mathrm{H} \alpha$

SNR shock is considered as the best candidate of CR origin.

## Shock Structure with CRs

## no CR

pobrotons 1 a
background plasha

## Shock Structure with CRs



CRs excite plasma waves and/or heat up the background plasma.

## If CRs are more

 efficiently accelerated, ...
## heated/turbulent

shock ( $z=0$ )

## Shock Structure with CRs

## no CR

Velocity $\uparrow$
CR pressure becomes
 considerably large.
The upstream plasma is decelerated.

$$
p=1+3 u_{2} /\left(u_{0}-u_{2}\right)
$$

The spectrum is determined by CRs themselves!
shock ( $z=0$ )

## Shock Structure with CRs

 never been observed in SNR.$\square$ The "heating" precursor can also be formed by photoionization (e.g. Ghavamian+00, Medina+14).
$\square$ The "velocity modification" would be firm evidence of the modification of shocks.

## Cosmic-Ray Modified Shock (CRMS)





Terasawa 2006
~20 \% of shock energy (flux) converted to non-thermal particles

- In situ. observation Solar wind
$\mathrm{V}_{\mathrm{sh}} \sim 100 \mathrm{~km} / \mathrm{s}$
$\mathrm{E}_{\mathrm{CR}} \sim 10 \mathrm{keV}-\mathrm{MeV}$
$B \quad \sim 10 \mu G\left(M_{A} \sim 5\right)$
Age ~day
Young SNR
$\mathrm{V}_{\mathrm{sh}}>1000 \mathrm{~km} / \mathrm{s}$
$\mathrm{E}_{\mathrm{CR}} \sim 1 \mathrm{GeV}-3 \mathrm{PeV}$ ?
B $\sim 1-100 \mu \mathrm{G}$ ?
( $\mathrm{M}_{\mathrm{A}} \sim 1-100$ ?)
Age $\sim 100-1000 \mathrm{yr}$.


## Cosmic-Ray Modified Shock (CRMS)





## Cosmic-Ray Modified Shock (CRMS)



We must examine such modification of plasma located at a distance of kpc scale. Challenging!


## Supernova Remnant (SNR)


$\gamma$-ray: electron or proton?

X-ray:~TeV CR electrons Supernova ejecta

Ha: useful tracer of shock condition \& physics.

H $\alpha$ emissions reflect a plasma condition around the shock (e.g. Raymond 91 for review). Red: H $\alpha$ best candidate of CR origin.

## H $\alpha$ emission from upstream



## H $\alpha$ emission from upstream



## H $\alpha$ emission from upstream



## Ha emission from upstream

$\mathrm{H}_{\mathrm{H}} \quad$ noCR Velocity -

H passes through the shock front $\stackrel{\downarrow}{\mathbf{H}}+\mathrm{p} / \mathrm{e} \rightarrow \underline{\mathrm{H}^{*}}+\mathrm{p} / \mathrm{e}$ emits Ly $\beta$
$\mathrm{p} / \mathrm{e}$ suffers the shock compression

## Ha emission from upstream

$\mathrm{H}_{\mathrm{H}} \quad$ no CR Velocity $\uparrow$

H passes through the shock front $\stackrel{\downarrow}{\mathbf{H}}+\mathrm{p} / \mathrm{e} \rightarrow \underline{\mathrm{H}^{*}+\mathrm{p} / \mathrm{e}}$ $\uparrow$ emits Ly $\beta$
$\mathrm{p} / \mathrm{e}$ suffers the shock compression

## H $\alpha$ emission from upstream



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$\mathrm{p} / \mathrm{e}$ suffers the shock compression

## H $\alpha$ emission from upstream

Velocity $\uparrow$

shock (z=0)

## H $\alpha$ emission from upstream



If we observe from the $y$-direction, the polarization angle is ...

shock ( $z=0$ )

## Polarization angle for $\mathrm{Ly} \beta \rightarrow \mathrm{H} \alpha$



## H $\alpha$ emission from upstream

Velocity $\uparrow$

shock (z=0)

## H $\alpha$ emission from upstream



The upstream $\mathrm{H} \alpha$ comes from:

1. Raman scattering resulting in Ly $\beta$ to $\mathrm{H} \alpha$
2. Collisional excitation in the upstream region.
shock ( $z=0$ )

## H $\alpha$ emission from upstream



## H $\alpha$ emission from upstream



## H $\alpha$ emission from upstream



H passes through the shock front $\stackrel{\downarrow}{\mathbf{H}}+\mathrm{p} / \mathrm{e} \rightarrow \underline{\mathrm{H}^{*}+\mathrm{p} / \mathrm{e}}$ emits Ly $\beta$
p/e suffers the shock compression
fully ionized shock ( $z=0$ )

## H $\alpha$ emission from upstream



## H $\alpha$ emission from upstream



Collisional excitation in the heating precursor can also yield Ly $\beta$ perpendicularly to the modification.
fully ionized
shock ( $z=0$ )

## Polarization angle for $\mathrm{Ly} \beta \rightarrow \mathrm{H} \alpha$



## H $\alpha$ emission from upstream

| Polarization direction (definition of Stokes $Q$ ): |
| :--- |
| no modification $\rightarrow$ parallel $\rightarrow Q<0$ |
| Modified $\rightarrow$ perpendicular $\rightarrow Q>0$ |



## H $\alpha$ emission from upstream

R Radiation line transfer \& atomic population with polarized light $\rightarrow$ quite complex!
$\square$ We make simplifications:

1. Omitting the polarization in atomic population calculations (SJ \& Laming 19a). Stokes $I$ is OK.
2. Completely unpolarized Ly $\beta$ is supposed.
3. For the Stokes $Q$, the $3 \mathrm{p}_{3 / 2}$ state only results from the radiative excitation in the upstream region.

See, SJ \& Laming 19b for details

## Model set up for the shock


precursor front $\left(z=z_{\text {pre }}\right) \quad$ shock/subshock $(z=0)$
Downstream value is given by usual Rankine-Hugoniot relations with $T_{\mathrm{e}}=0.1 T_{\mathrm{p}}$ (incomplete ion-electron temperature equilibrium).

## Model set up for the shock



We solve 3 cases:

1. No precursor
2. electron heating precursor, but no modification
3. electron heating precursor with decelerated protons, but no proton heating (Cosmic-Ray Modified Shock)

precursor front $\left(z=z_{\text {pre }}\right) \quad$ shock/subshock $(z=0)$
Downstream value is given by usual Rankine-Hugoniot relations with $T_{\mathrm{e}}=0.1 T_{\mathrm{p}}$ (incomplete ion-electron temperature equilibrium).

## Line Transfer Model



## Parameters:

(1) Shock velocity $V_{\text {sh }}$
(2) Upstream number density $n_{\text {tot }, 0}$
(3) proton fraction $\chi_{0}$
(4) Upstream electron temp
(5) Downstream electron temp $T_{\mathrm{e}}=\beta T_{\text {down }}$

Pure hydrogen plasma.
We solve the excited states up to 4 f .
(SJ \& Laming 19a)

## Results: Ionization Structure of H



## Results: Radiative vs. Collisional



## Results: Polarization of $\mathrm{H} \alpha$

The sign of degree indicates the polarization angle (Stokes $Q$ ).


## Results: Polarization of $\mathrm{H} \alpha$

Polarization degree


Surface brightness


The polarization direction can respond whether the shock is modified.

The degree of a few per cent is measurable (Sparks+ 15).

## Discovery of polarized $\mathrm{H} \alpha$ emission

@ bright filament of SN 1006 (Sparks+ 15)



$>$ Linear Polarization
$>$ Polarization angle: perpendicular to the shock
$>$ Degree: $2.0 \pm 0.4$ \%

## Discovery of polarized $\mathrm{H} \alpha$ emission

@ bright filament of SN 1006 (Sparks+ 15) 6000 5000
$\checkmark$ Our calculation is consistent with the observation.
$\checkmark$ Polarized H $\alpha$ has been reported by Sparks +15 , but this is not spatially resolved...
$\checkmark$ Further observations of $\mathrm{H} \alpha$ polarimetry will bring new insights to particle acceleration!

|  | $>$ Linear Polarization |
| ---: | :--- |
|  | $>$ |
|  | Polarization angle: |
|  | perpendicular to the shock |

## Polarization of $\mathrm{H} \alpha$ vs others

Polarization degree

$H \beta / H \alpha$


Surface brightness


## Only polarization can catch the velocity modification!

## Summary

$\square$ Cosmic-Ray Modified Shocks (CRMSs) are one of essential prediction of the diffusive shock acceleration.
$\square$ We must examine a velocity modification of plasma with $\sim 10$ \% level around the SNR shock.
$\square$ The polarization direction of $\mathrm{H} \alpha$ responds sensitively whether the shock is modified.

## SN 1006



## SN 1006





$$
\begin{aligned}
L_{\pi^{0}} & \sim 0.1 \frac{\eta \rho V_{\mathrm{sh}}^{2}}{2} R^{3} n \sigma c \\
& \sim 10^{33} \eta \mathrm{erg} / \mathrm{s} \times\left(\frac{R}{3 \mathrm{pc}}\right)^{3}\left(\frac{V_{\mathrm{sh}}}{0.01 c}\right)^{2}\left(\frac{n}{0.3 \mathrm{~cm}^{-3}}\right)^{2}
\end{aligned}
$$

# Balmer Line Emissions from Collisionless Shocks 

Winkler+14
Supernova Remnants (SNRs)


Figures from Morlino+15
Pulsar Wind Nebulae

Balmer line emissions (especially $\mathrm{H} \alpha$ ) are ubiquitously seen in collisionless shocks propagating into the ISM.

## Balmer Line Emissions from Collisionless Shocks



# Balmer Line Emissions from Collisionless Shocks 

$\checkmark$ Emission Mechanism (e.g. Chevalier+80) upstream downstream • The collisionless shock is
 formed by the interaction between charged particles and plasma waves.

- The neutral particles (e.g. hydrogen atoms) are not affected.

SNR shock
Charged particles $\rightarrow$ shock heating
Hydrogen atoms $\rightarrow$ no dissipation

# Balmer Line Emissions from Collisionless Shocks 

$\checkmark$ Emission Mechanism (e.g. Chevalier+80)


SNR shock
Charged particles $\rightarrow$ shock heating Hydrogen atoms $\rightarrow$ no dissipation

- Collisional Excitation

$$
\mathrm{H}+\mathrm{p}(\text { or } \mathrm{e}) \rightarrow \underline{\mathrm{H}}^{*}+\mathrm{p}(\text { or } \mathrm{e})
$$ Emits "narrow" comp.

$\square$ Charge Transfer
$\mathrm{H}+\mathrm{p} \rightarrow \mathrm{p}+\underline{\mathrm{H}^{*}}$
Emits "broad" comp.
The "broad" component reflects the downstream temperature of protons.

